

Surfaces, Interfaces, Spins

– The age of Interface!

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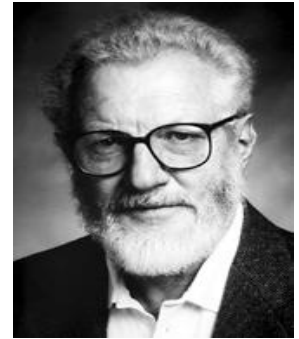
Why interface is important?

“The interface is the device”

- Herbert Kroemer

Interface is the platform for creating exotic electronic states:

- **Semiconductor heterostructure interface:** (*electron-electron interaction*)
high mobility electron gas, fractional charge and statistics, topological order etc.
- **Oxide heterostructure interface:** (*correlated electrons*)
high mobility electron gas, ferroelectricity, high temperature superconductivity etc.
- **Topological insulator surface:** (*spin-orbit coupling*)
topological surface states
- **Combined interactions:** ferromagnetic, superconductivity, spin-orbit etc.
→ exotic electronic state: Majorana fermion



Herbert Kroemer

Nobel Prize in Physics (2000)

"for developing semiconductor heterostructures used in high-speed- and opto-electronics"
– Nobelprize.org

***Postdocs:* Ferhat Katmis, Peng Wei, CuiZu Chang
Yunbo Ou , Juanpedro Cascales**

Visiting Scientist: Yota Takamura

Collaborators

mK transport - W.W. Zhao, J. Jain and M. Chan - Penn State

PNR - Valeria Lauter - Oakridge Natl Lab

SQUID - Don Heiman - Northeastern U

XAS, XMCD - John Freeland - Argonne Natl Lab

XTEM - Biswarup Satpathy - Saha Inst. Nucl. Physics, India

TIG/TI - Chi Tang, Jing Shi - U of California Riverside

Theory - Kyungmin Lee, Nandini Trivedi - Ohio State U.

- Spins - Spintronics: unconventional
- Spins - Superconductors
- Spins - Quantum coherent systems

The common platform

- surfaces and interfaces!!

Towards Molecular Spintronics

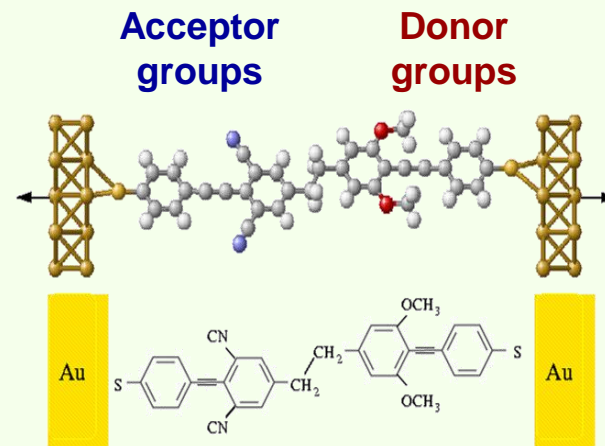
Why single molecules?

- ❖ Smallest well defined electronic units
- ❖ Chemically synthesize precisely designed properties, with atomic precision - tune tailored molecules !

..... Unlimited possibilities

- ❖ Easily scalable

Add: electron spin degree of freedom for spintronic functions within single molecular building block



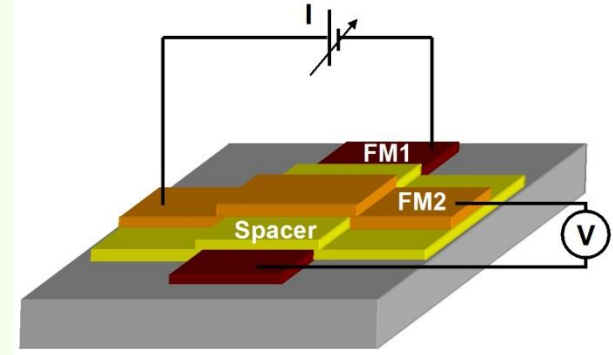
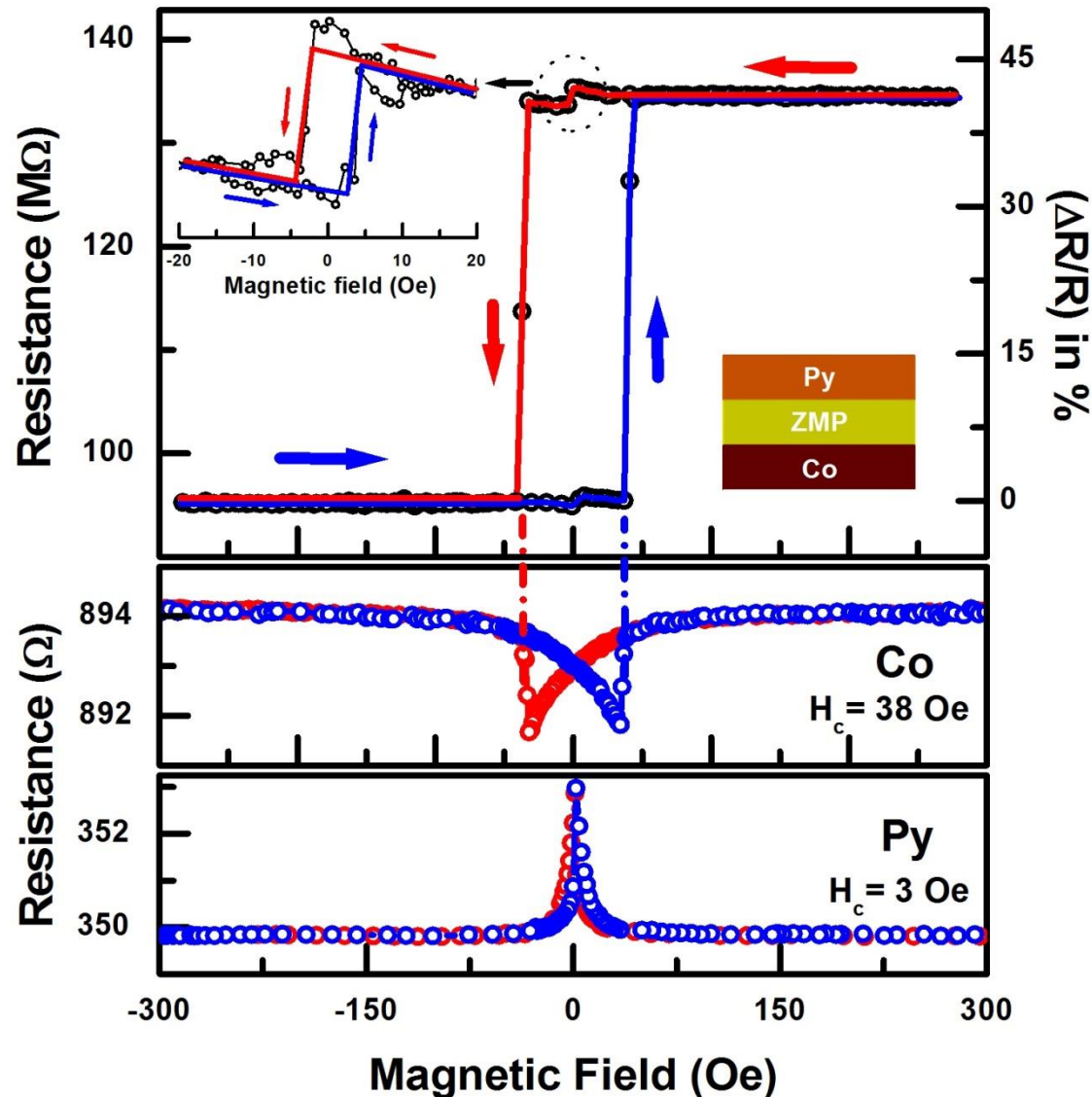
Molecular rectifier
Aviram and Ratner, (1974).

➡ Towards single molecule information storage and sensing

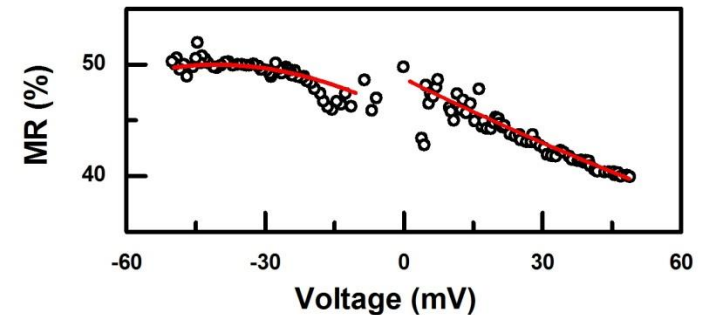
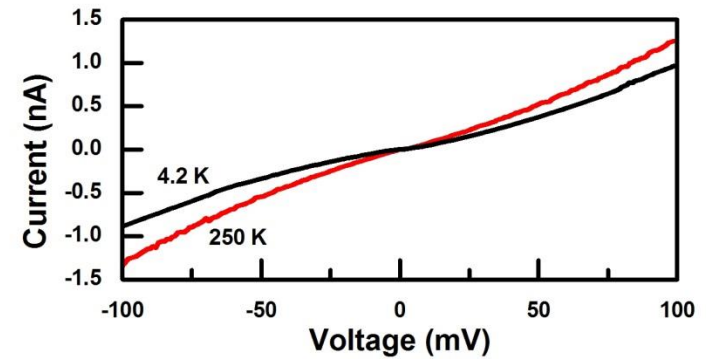
We are at the beginning → Need of basic research!

Magnetoresistance in Co/ZMP (40 nm)/Py spin valve

$V_{dc} = 25$ mV and $T =$ at 4.2K

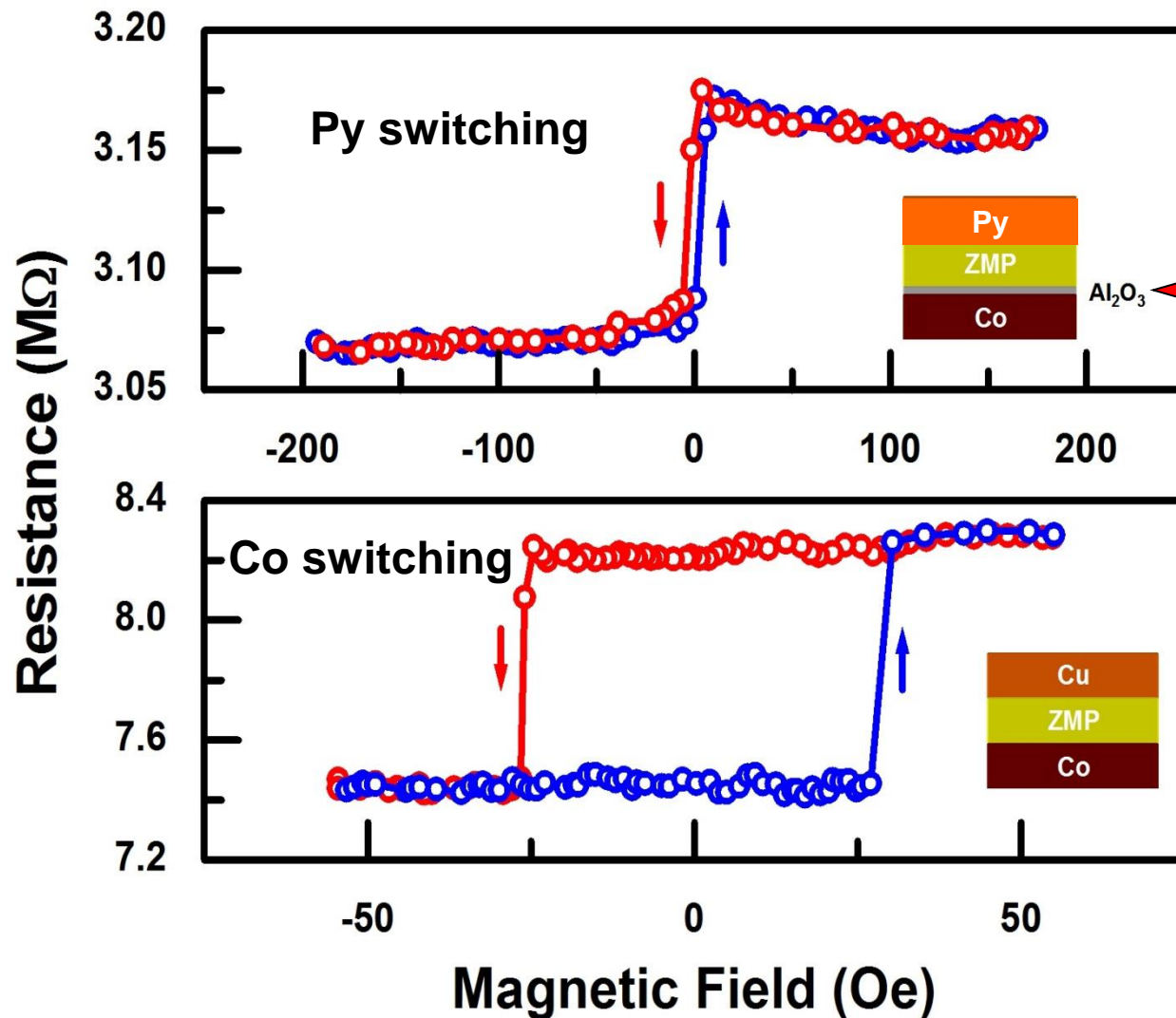


Non-linear characteristics



Interface effect: K. V. Raman JSM, Nature 493, 509 (2013)

Interface Effect: Confirmation!



No induced FM or antiparallel spin state at the bottom interface with 0.3nm Al_2O_3 inbetween

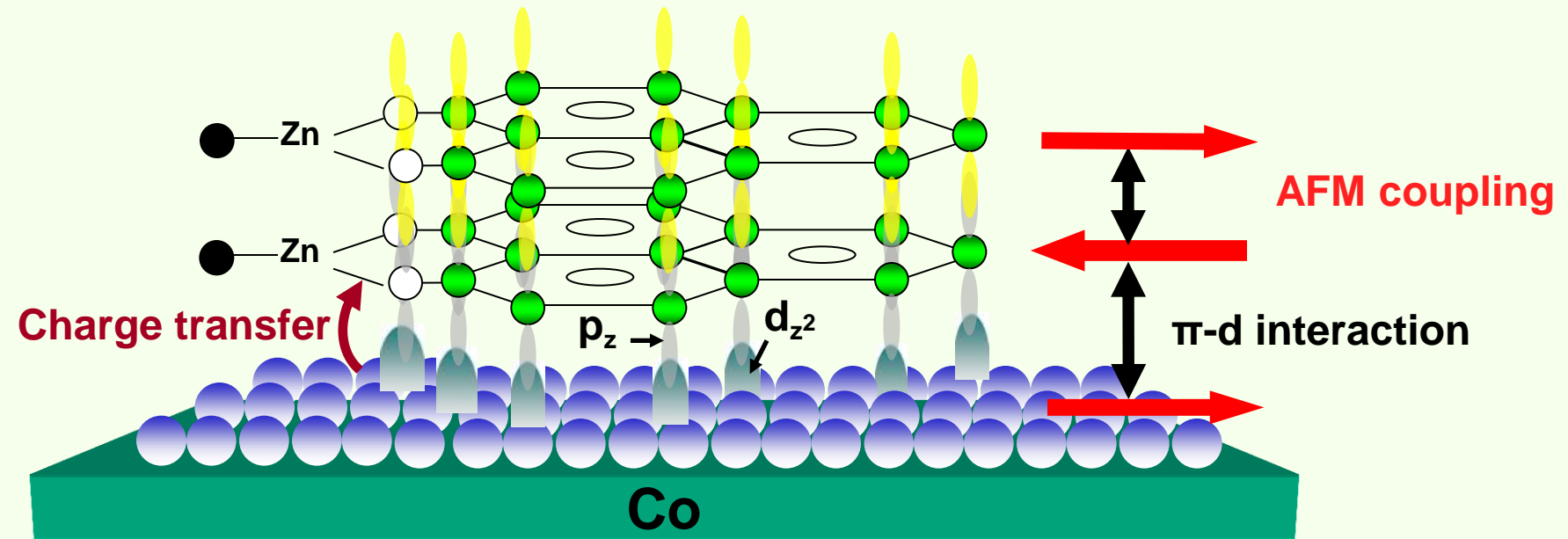
No magnetism or switching at the top interface with Cu

No MR in Cu/ZMP/Cu

A 'pinned magnetic layer' exists at the interface

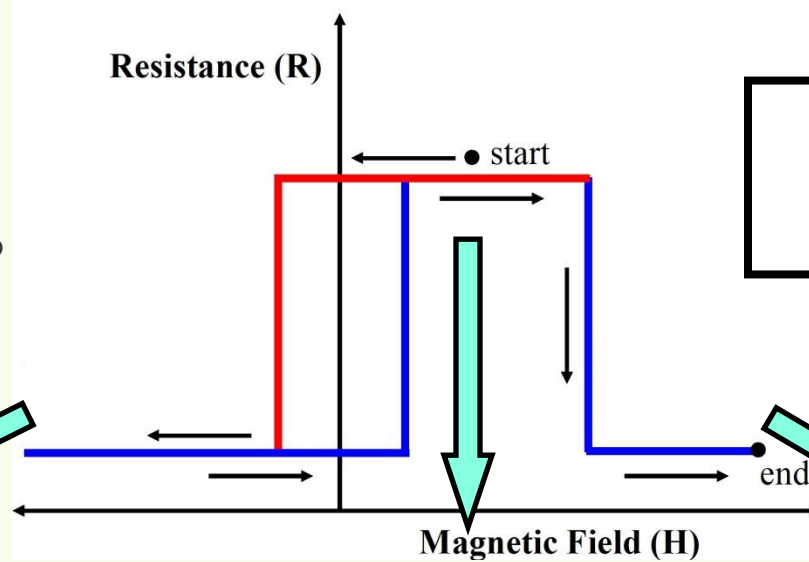
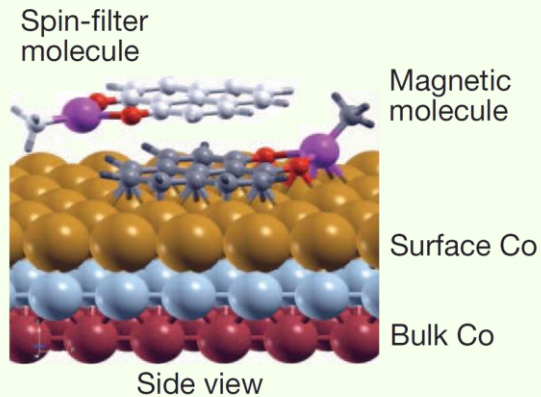
Interface phenomena

Zinc methyl phenalenyl molecules over Co surface



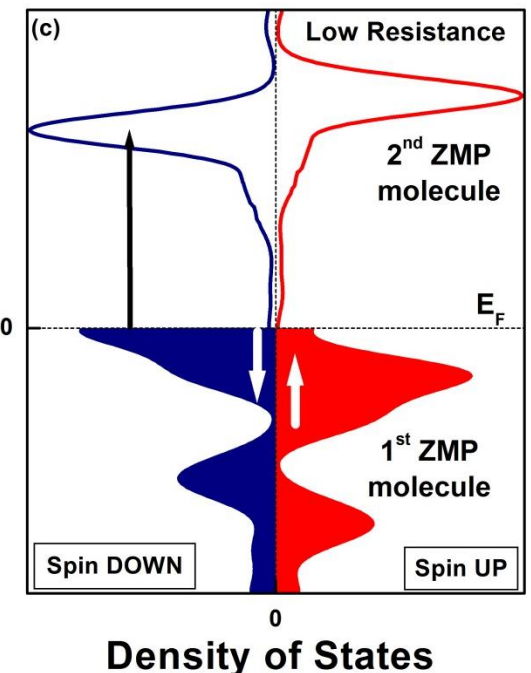
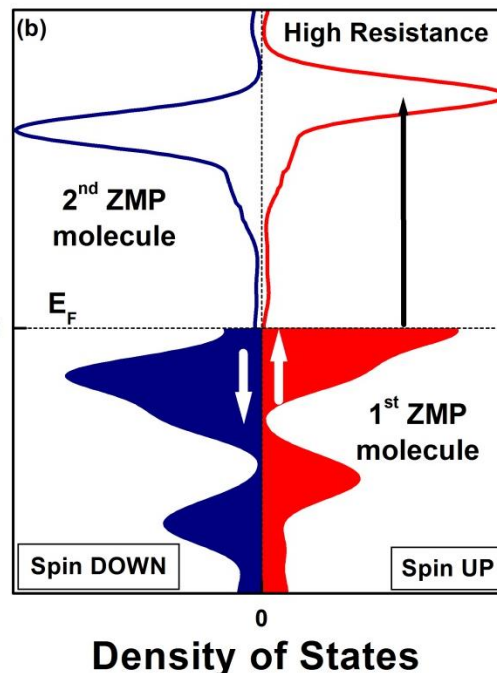
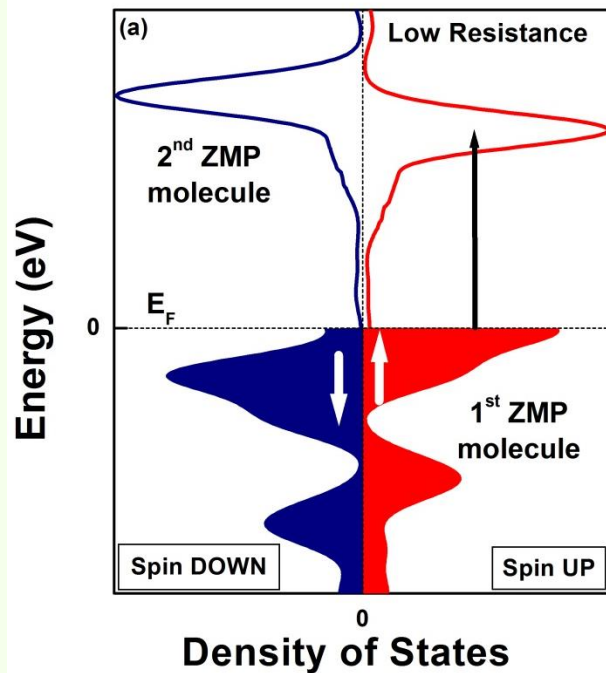
Origin of IMR effect

Raman et al. Nature 493, 509 (2013)



**Spin-filter
injection/tunneling
response**

**Theory by:
N. Atodiresei**



Towards molecular spintronics

- planar organic molecules (ZMP)
 - magnetic & spin-filtering - room temperature
 - Molecules can store information as a classical bits i.e. '0' or '1' state
- Storage density may reach ~ 1000 TBytes/sq inch

Technological advantages ...

Only one FM needed

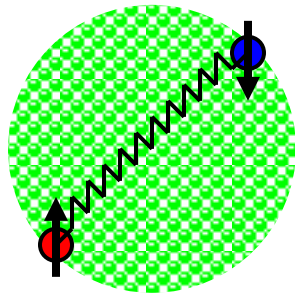
Molecular 'bit' - integrated storage-sensing system

Molecules are ideal quantum dot systems

- fabrication of scalable quantum memory registers

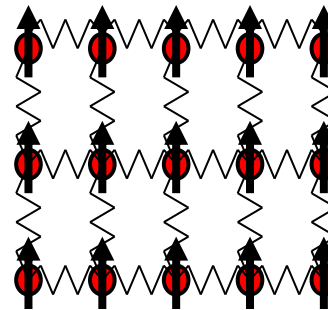
Coexistence of Superconductivity *and* Ferromagnetism !!

Superconductivity



$$E_{\text{BCS}} \sim 10^{-3} \text{ eV}$$

Ferromagnetism



$$E_{\text{EX}} \sim 1 \text{ eV}$$

SC and FM are competing spin ordering mechanisms

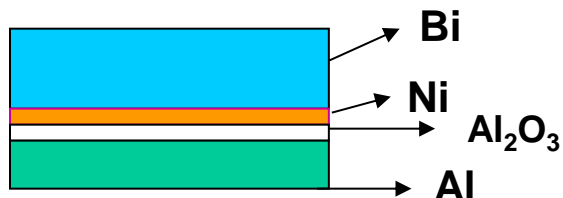
*An unexpected observation
Ferromagnetic Ni is superconducting !!*

Ferromagnetic Ni is also superconducting...

Normally Bi not a superconductor and Ni is a FM

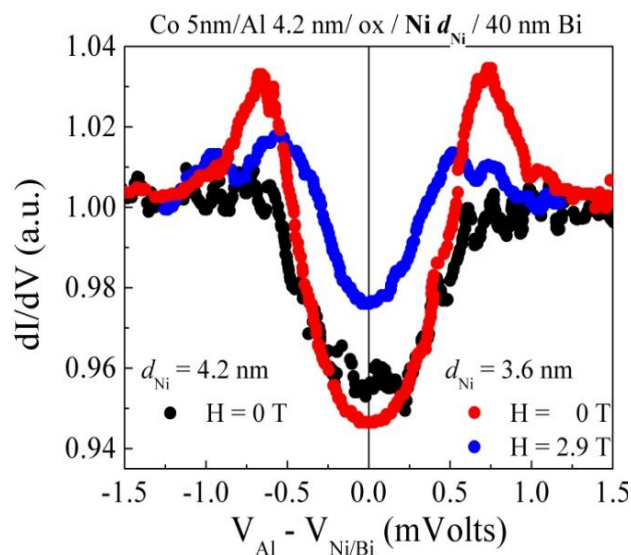
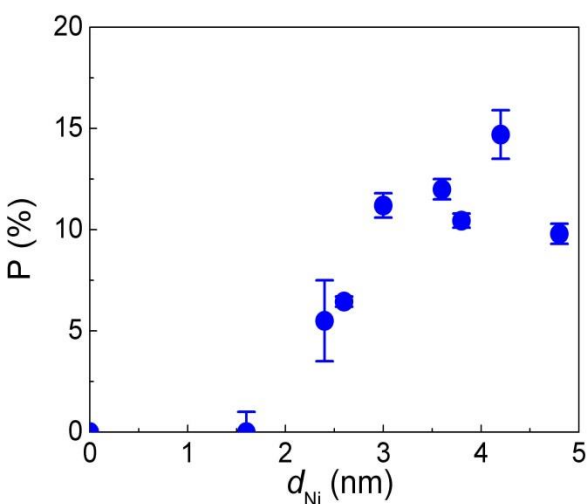
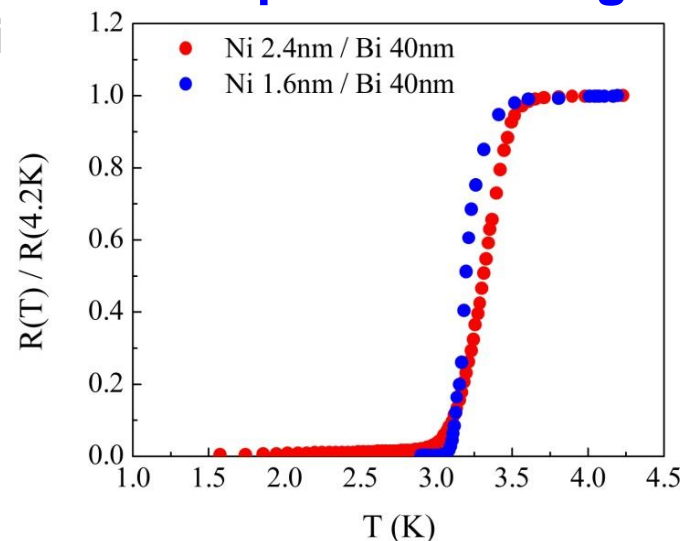
Combining the two – Ni/Bi bilayers

SC and FM !

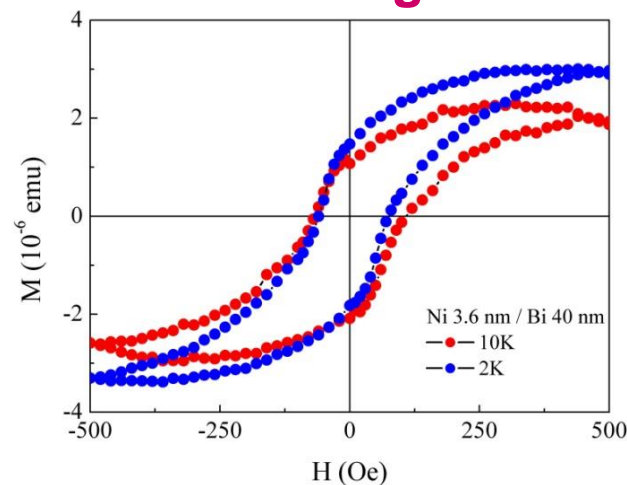


SC present in *whole* Ni/Bi stack

Superconducting !



Ferromagnetic !



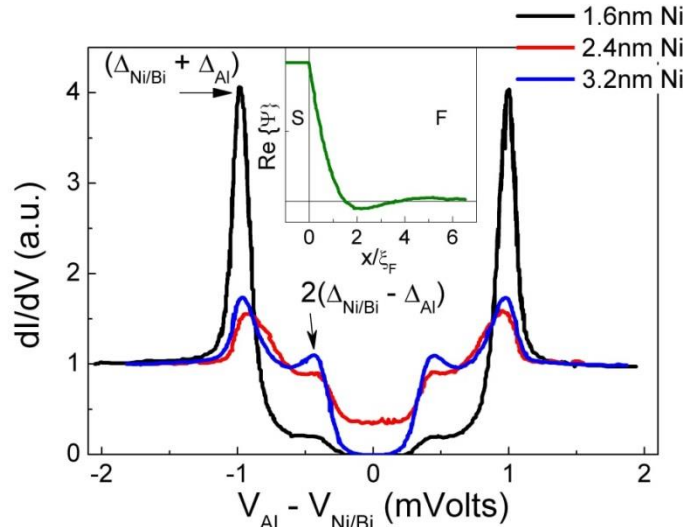
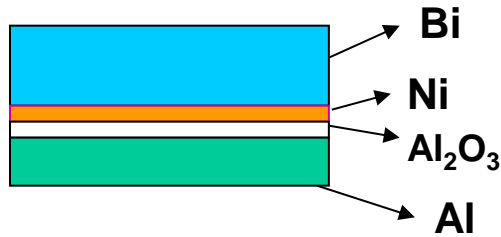
Superconductivity and ferromagnetism coexist in Ni !!

Spin polarized current observed in a triplet SC?!

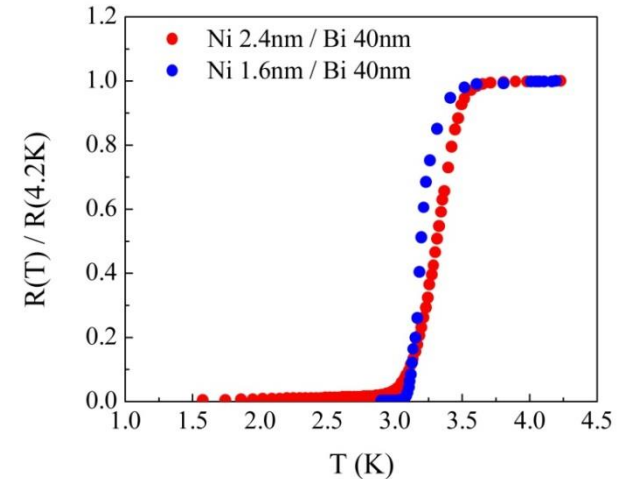
PRL 94, 037006 (2005)

Superconductivity and ferromagnetism coexist in Ni/Bi bilayers !!

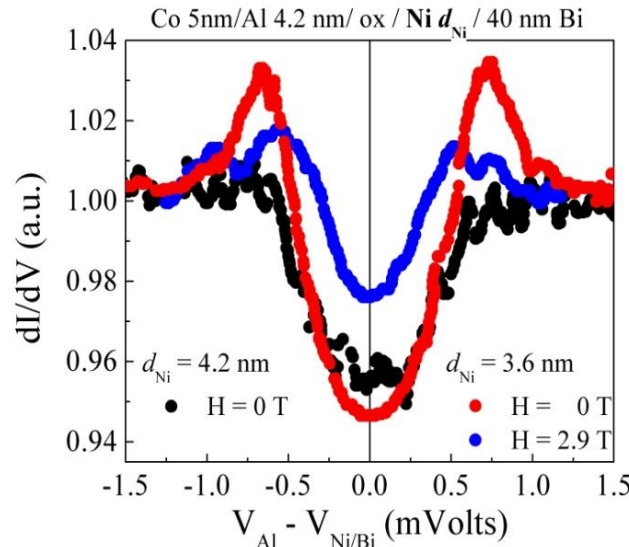
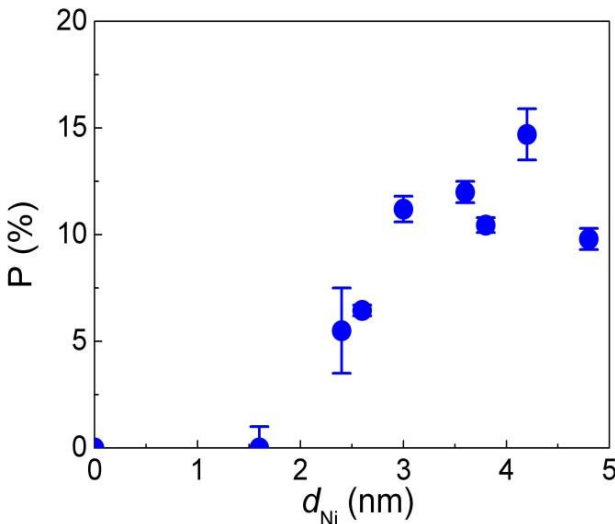
Normally Bi not a superconductor and Ni is a FM



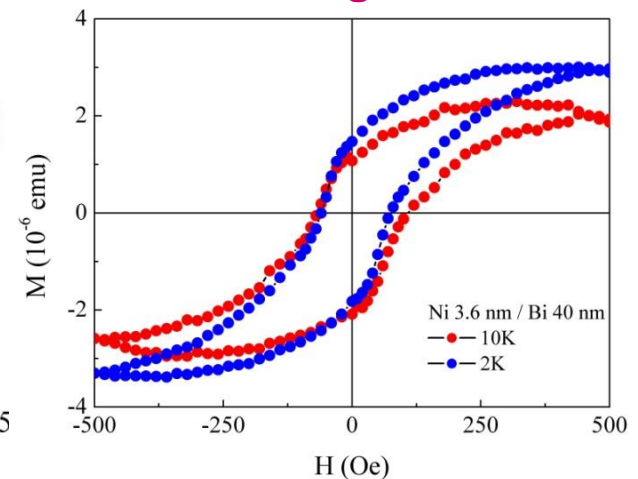
Superconducting



Spin polarized!



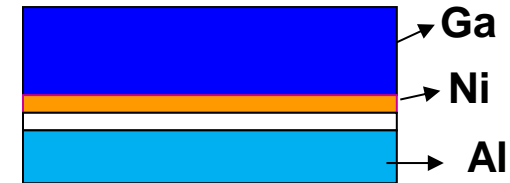
Ferromagnetic !



**Spin polarized current observed
- triplet paired SC?!**

PRL 94, 037006 (2005)

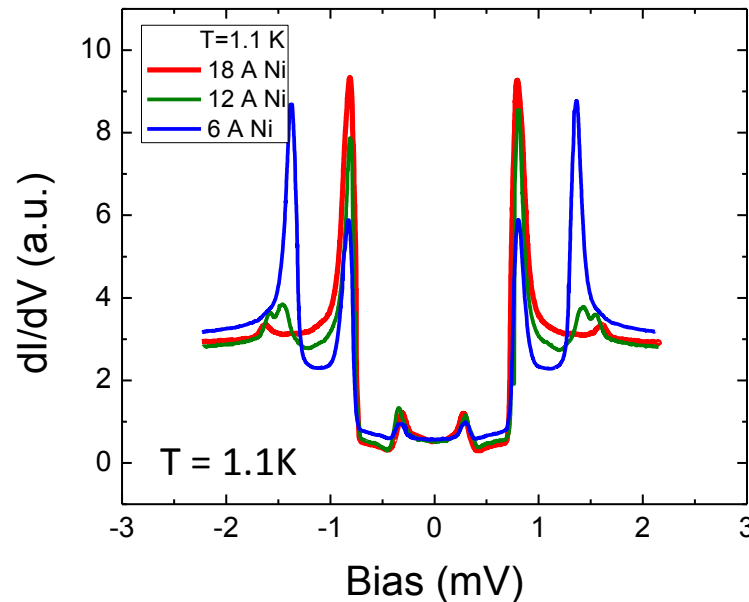
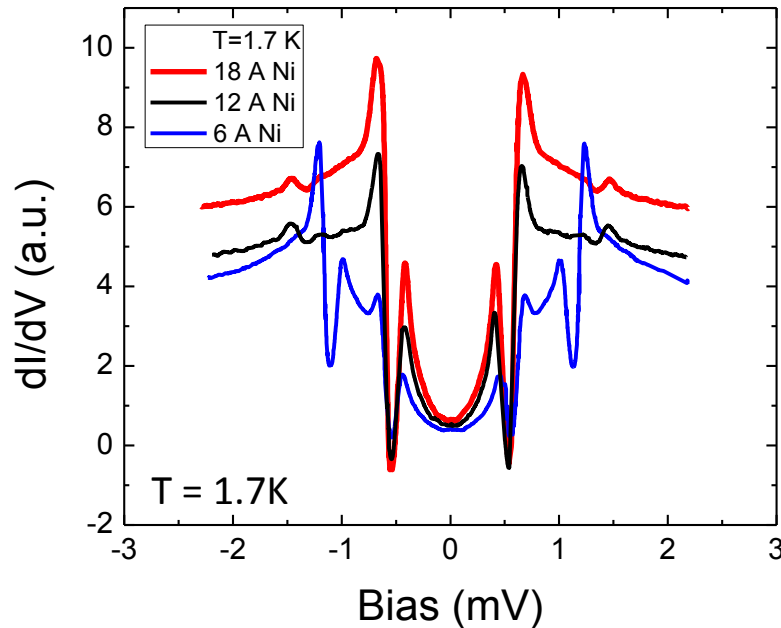
Bulk Ga $T_c = 1.08\text{K}$
 amorp Ga $T_c = 8.6\text{K}$, only stable $< 20\text{K}$



X Ni/600Å Ga film – superconducts with $T_c \sim 6.5\text{K}$

3-10057

S – I – S Tunneling: Three SC gaps !!



$T_c(\text{Al}) = 1.87\text{K}$
 At 1.7K
 $2\Delta_{\text{Al}} = 0.26\text{meV}$

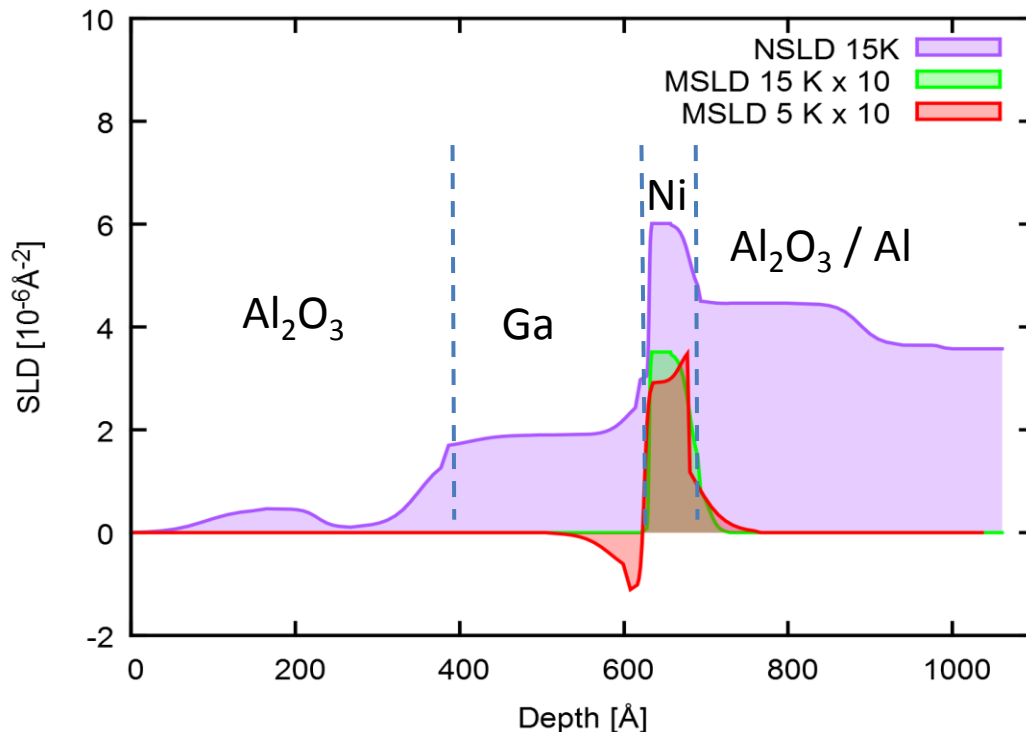
For 0.6nm Ni: $2\Delta_{\text{Surface}} = 1.14\text{meV}$; $2\Delta_{\text{Interface}} = 2.28\text{meV}$; $2\Delta_{\text{Bulk}} = 2.76\text{meV}$

For 2.4nm Ni: $2\Delta = 0.45\text{meV}$; $2\Delta = 1.15\text{meV}$; $2\Delta = 2.70\text{meV}$
 $2\Delta/kT_c \sim 4.5$ – strong coupled

Polarized Neutron Reflectivity results for Ni 56 A / Ga 600 A

At 5K Meissner state: field from the Ni side is screened at the Ga interface

M profile: positive M in Ni ; negative in Ga



Samples zero field cooled to 15 K,
H 1 kOe applied – measure PNR
ZFC to 5 K H = -150 Oe applied
H 1 kOe applied – measure PNR

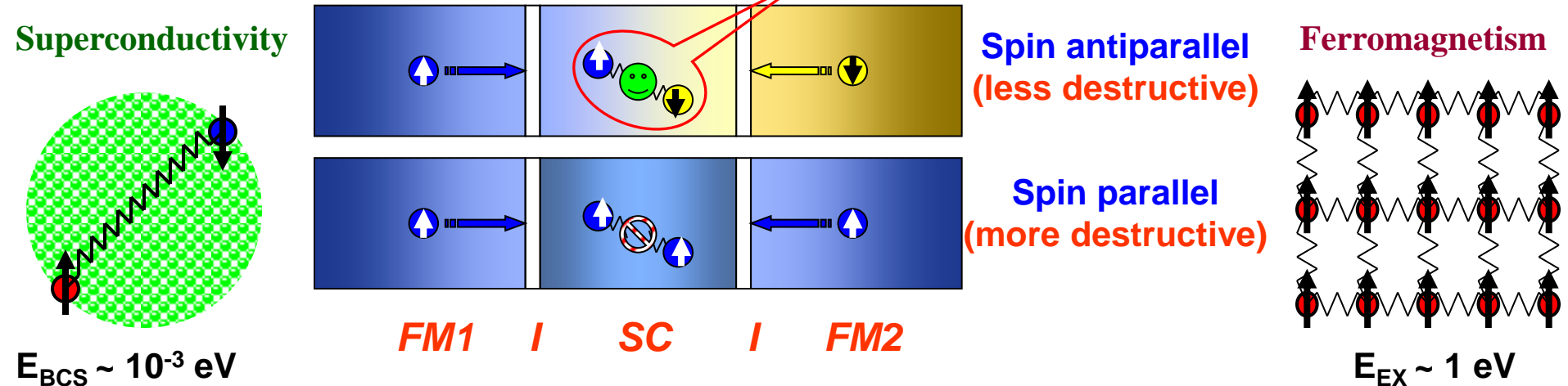
**40 and 56 A Ni show clear magnetic signature
adjacent to 600A Ga which is superconducting**

Tuning superconductivity with spin current ...

SC and FM are competing spin ordering mechanisms

GuoXing Miao

Cooper pairs



Large MR is expected between $\uparrow\uparrow$ and $\uparrow\downarrow$ states

Reduced transparency: 6 Fe / 1 Al_2O_3 / 4 Al / 1 Al_2O_3 / 6 Py > 1000% MR

Transparent interface: Epitaxial 6 Fe / 40 V / 6 Fe / 1 CoO – Infinite MR

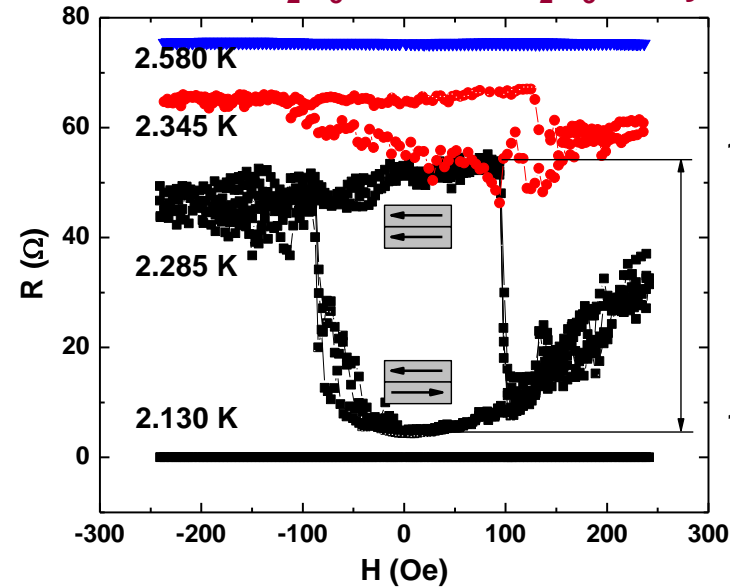
PRL 2007, 2008

Spin injection into superconductors

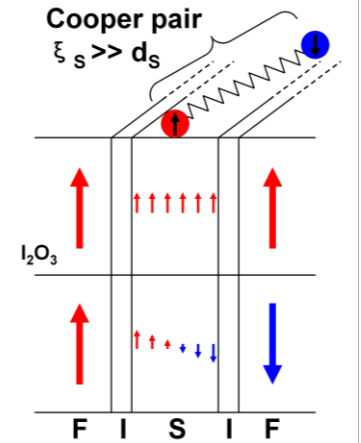
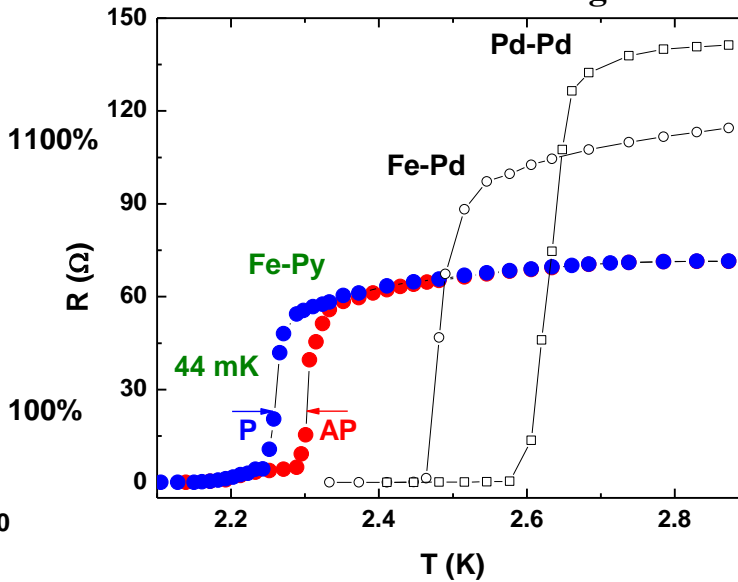
Over 1000% MR achieved!!

PRL 2007, 2008

6 Fe / 1 Al₂O₃ / 4 Al / 1 Al₂O₃ / 6 Py

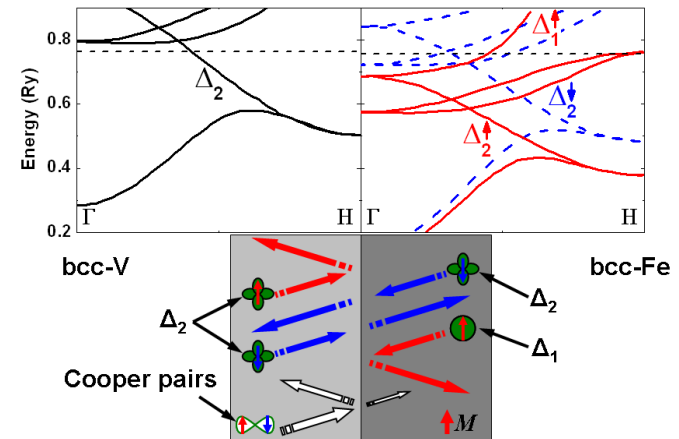
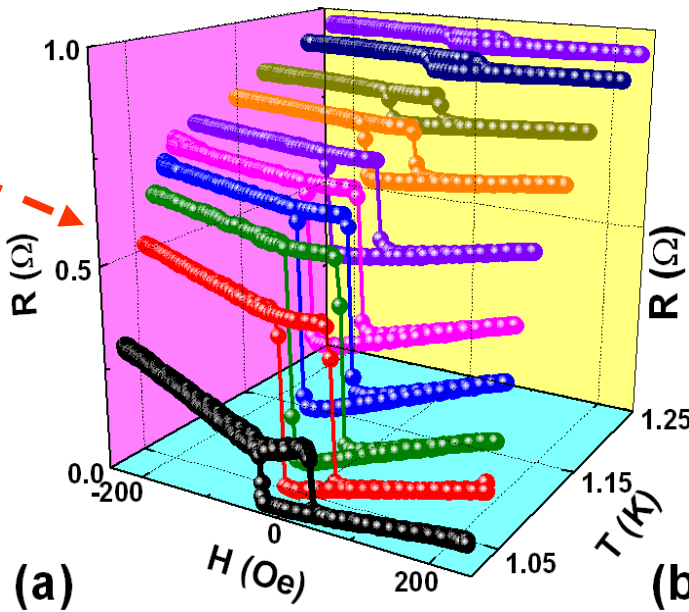
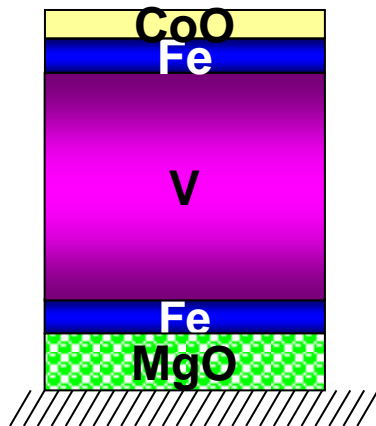


T_c shifts for P and AP alignment



BCC Heteroepitaxial Structure –

Band symmetry driven
Infinite MR achieved

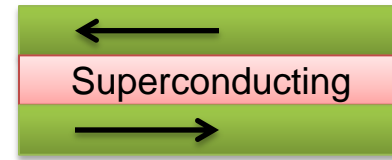
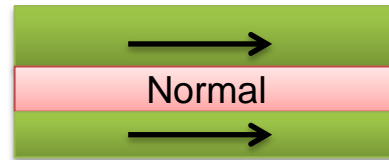


Conduction band in V is of Δ_2 symmetry
Absent in Fe spin majority channel.
Opaque interface for Cooper pairs

De Gennes' prediction – 50 years ago

de Gennes, *Phys. Lett.* **23**, 374 (1966)

Ferromagnetic Insulator
Superconductor
Ferromagnetic Insulator



Finite R



$R = 0$

Average exchange field

$$\bar{h} = 2|\Gamma|S \left(\frac{a}{d_s} \right) \cos \frac{\theta}{2}$$

From T_c change - exchange integral

$$\Gamma = 8.1 \text{ meV}$$

Average exchange field

$$\bar{h} \sim 13 \text{ meV}$$



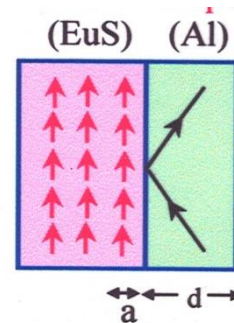
Shows perfect superconducting
spin switching effect

Equivalent external field $\sim 200 \text{ T}$

Internal exchange field – Proximity effect at a ferromagnet insulator / metal interface

Magnetic semiconductors EuS or EuO

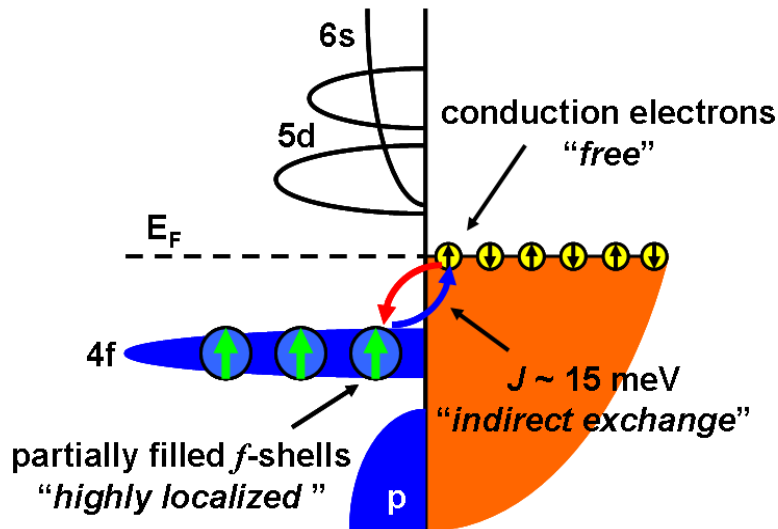
Exchange Interaction between EuS and normal metal



EuS

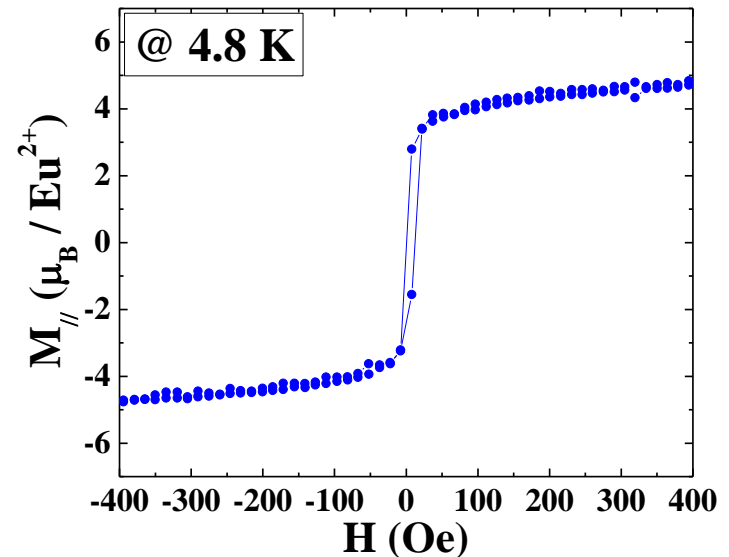
Band gap $E_g = 1.64$ eV

$\Delta E_{\text{ex}} = 0.36$



$$H_{EX} = -J (\vec{S}_{e^-} \cdot \vec{S}_{Eu^{2+}})$$

Magnetic Hysteresis 1 nm EuS thin film



Sarma 1963

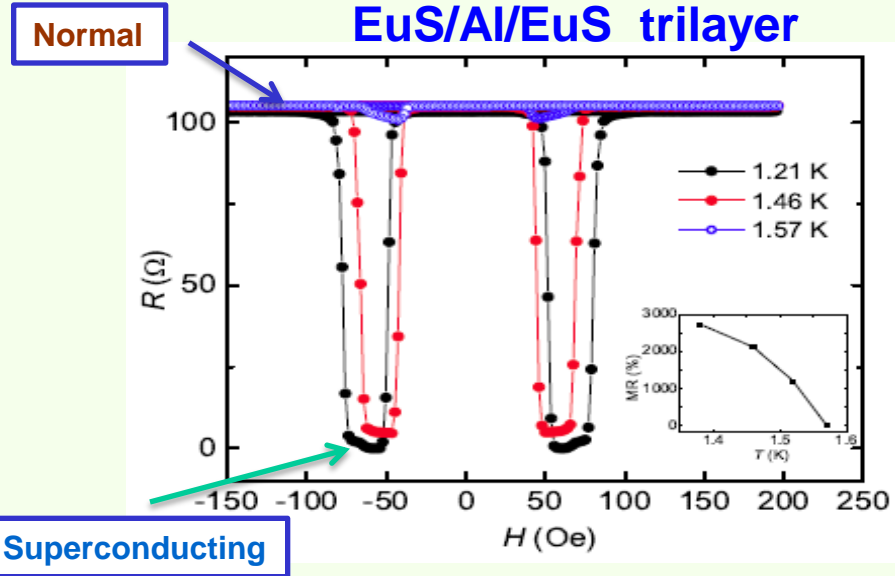
de Gennes, 1966

Tedrow et al PRL 1986

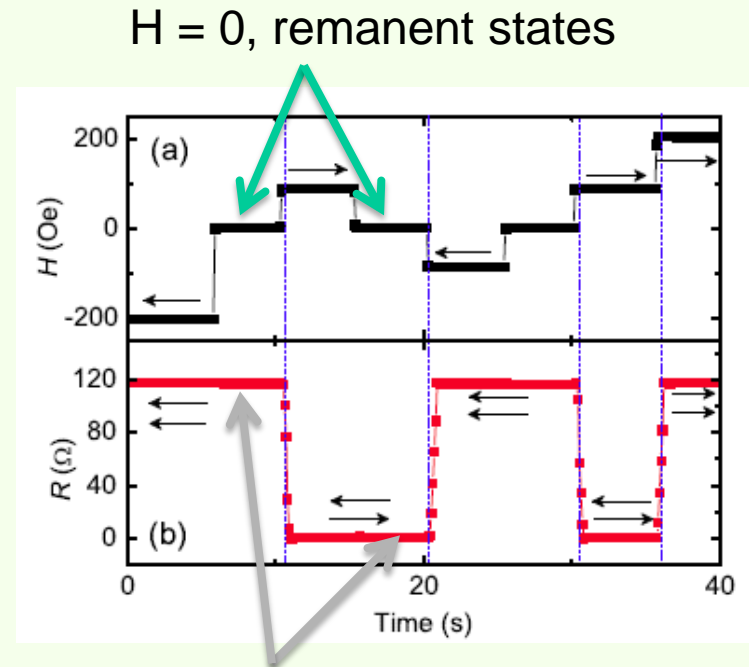
$$B^* \propto 1/d$$

EuS/Al/EuS superconducting spin switch

Li... JSM, *PRL* 2013



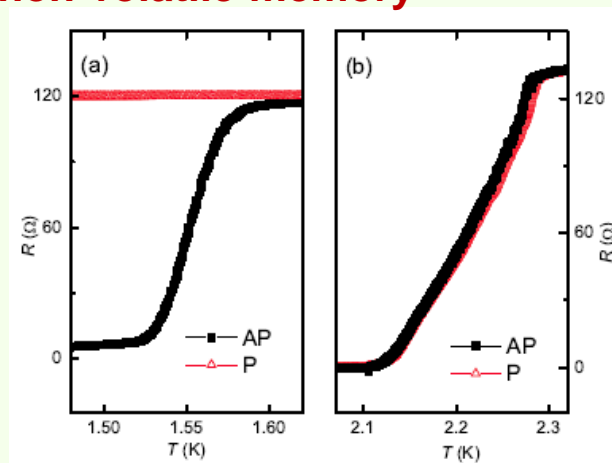
- Complete normal-superconducting transition
- **Infinite magnetoresistance (MR) ---**
Perfect switch, non-volatile memory



Two resistance states

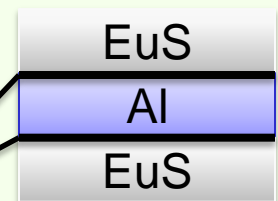
Supercond. transition

$$T_C^{AP} > T_C^P \longrightarrow$$



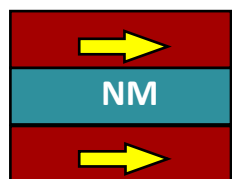
Suppresses proximity Exch.

0.3 nm Al_2O_3

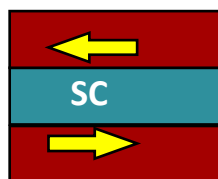


Superconducting Spin Switch realized at higher temperatures

Parallel

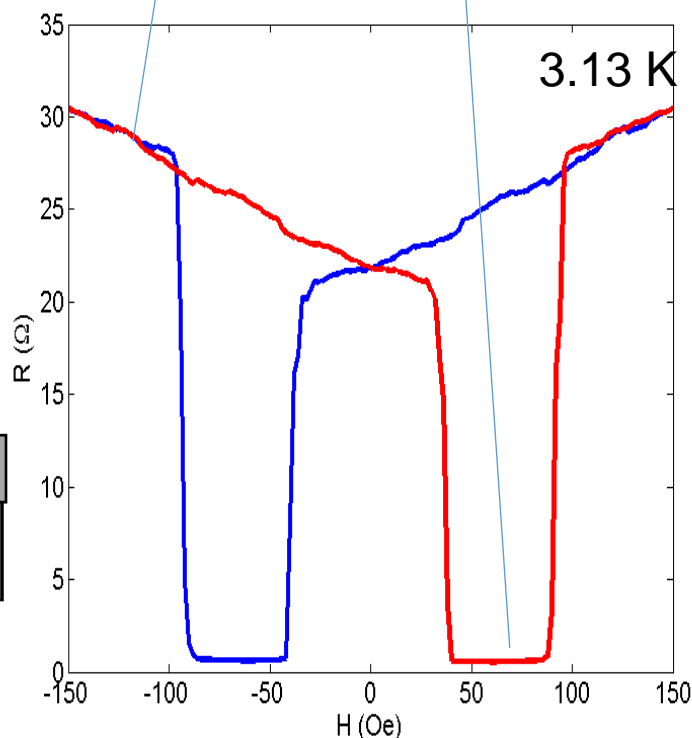
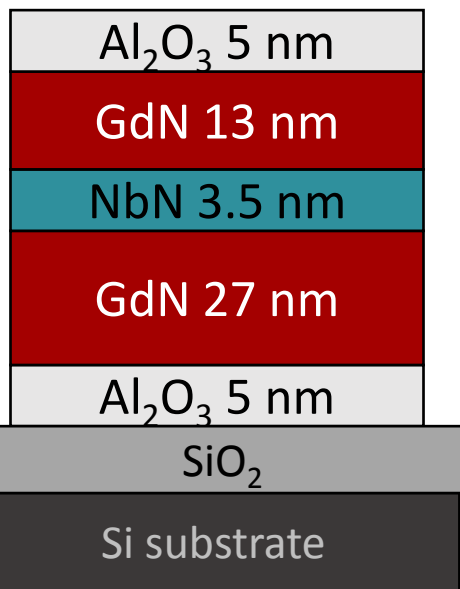


Anti-parallel

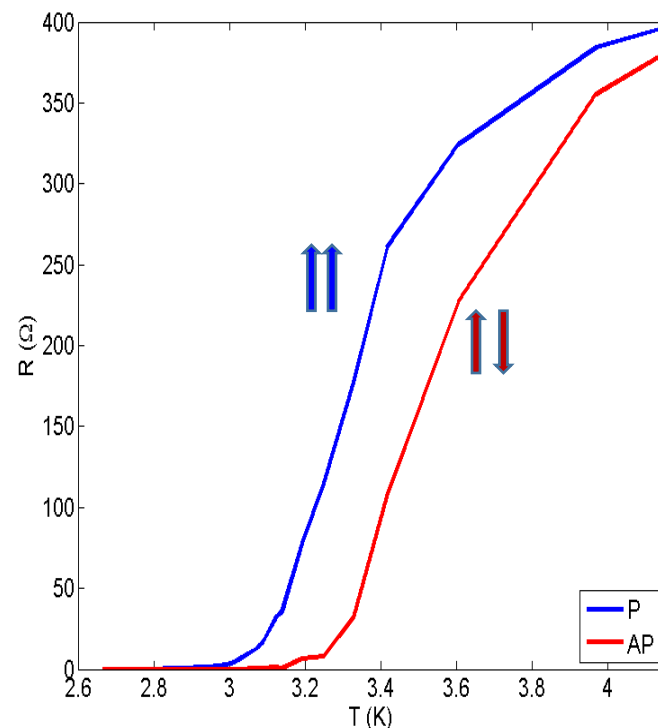


SC: Nb – 9.3K, NbN – 15-16K

FMI: GdN – ferromag. Tc 60K



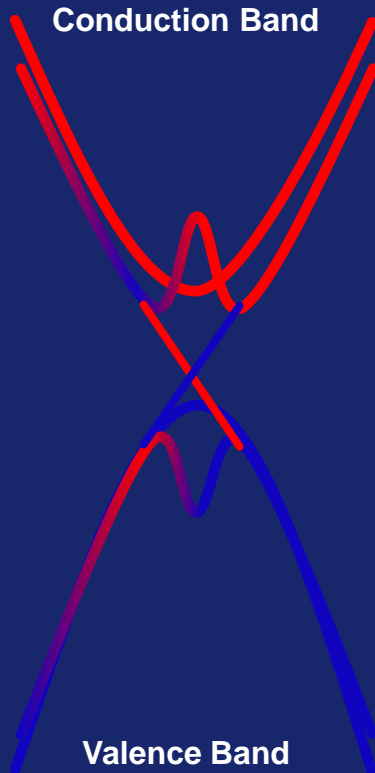
SC Resistive transition



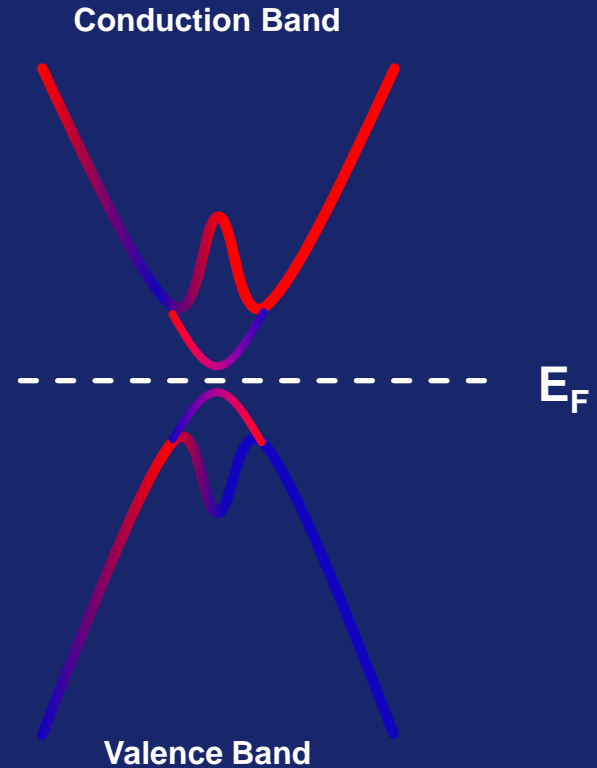
Perfect spin switching
Interface driven

Topological Insulators (TIs)

Topological Insulators

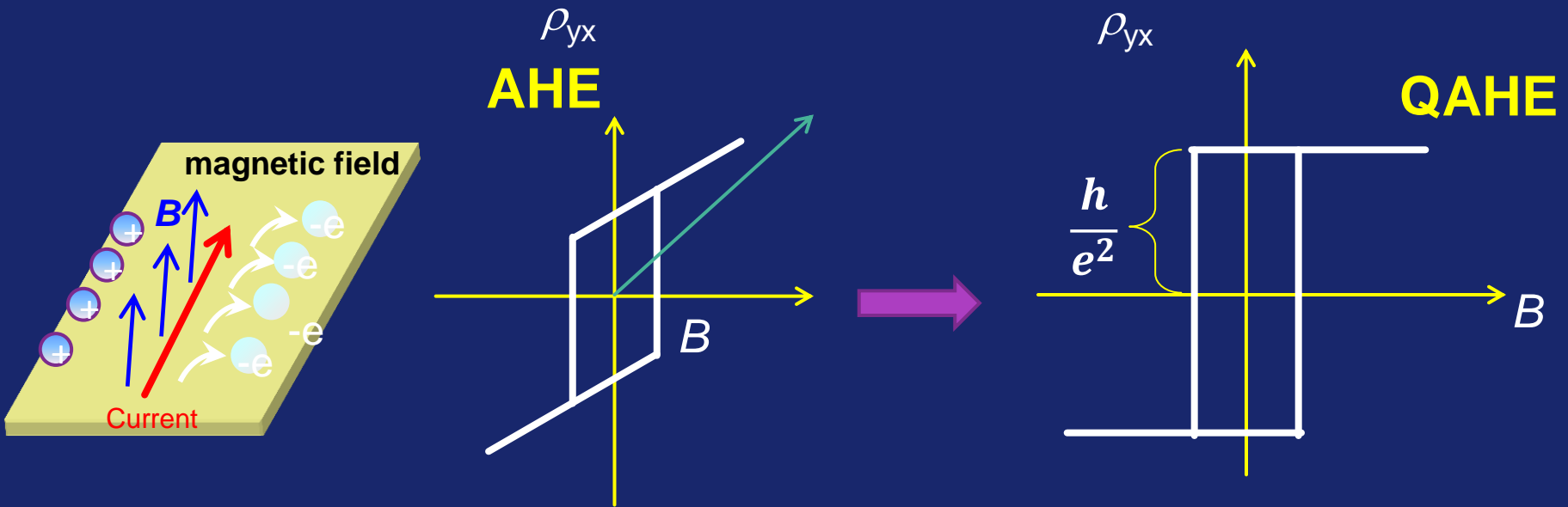


+ Ferromagnetic order



Quantum anomalous Hall (QAH) effect

Quantum anomalous Hall (QAH) effect



Hall effect at zero field
in a magnetic material

Quantum Hall effect at zero field

- Ferromagnetic insulator
- Topologically non-trivial electronic band structures



$$\rho_{yx} = \frac{h}{e^2} = 25.812 \text{ K}\Omega$$

$$\rho_{xx} = 0$$

Spin polarized ballistic chiral edge mode

- **TI thin films and heterostructures - High quality**

- Homogeneous composition
- Negligible defects
- High mobility and low carr. density

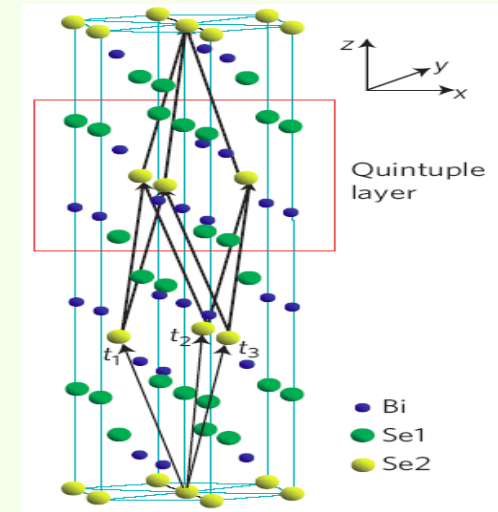
To reach QAH state

- **Create ferromagnetic TI by doping**
 - *where does the dopant go?*

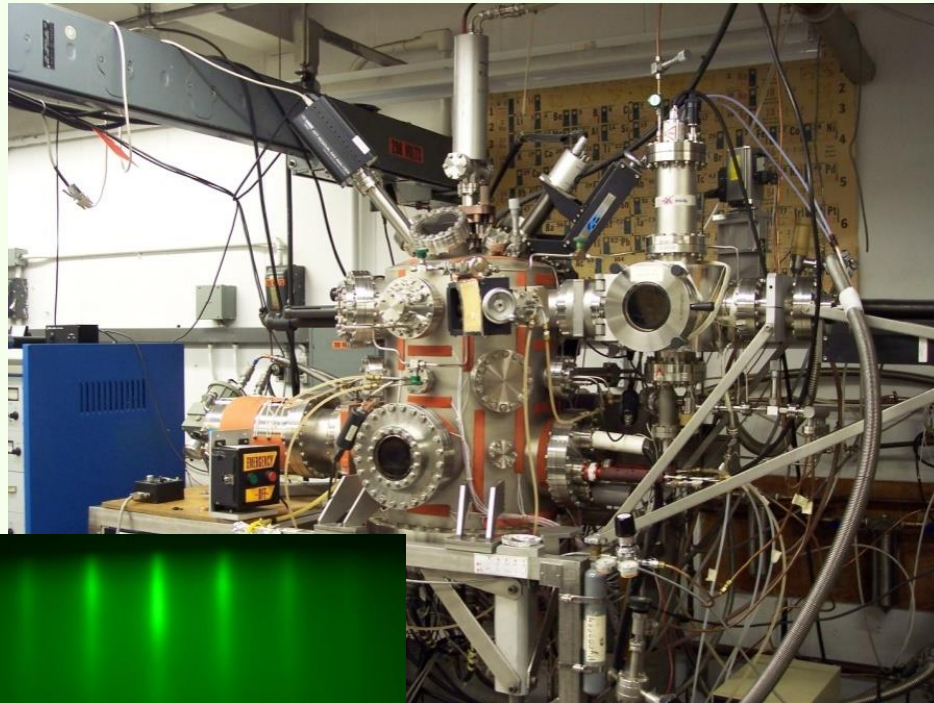
- **TI/FMI heterostructure**

Interface coupled magnetic behavior

Epitaxial growth (MBE) of TI layers

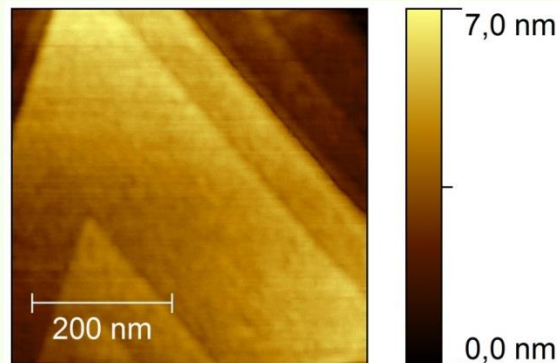


H. Zhang et al Nat Phys 5, 438 (2009)



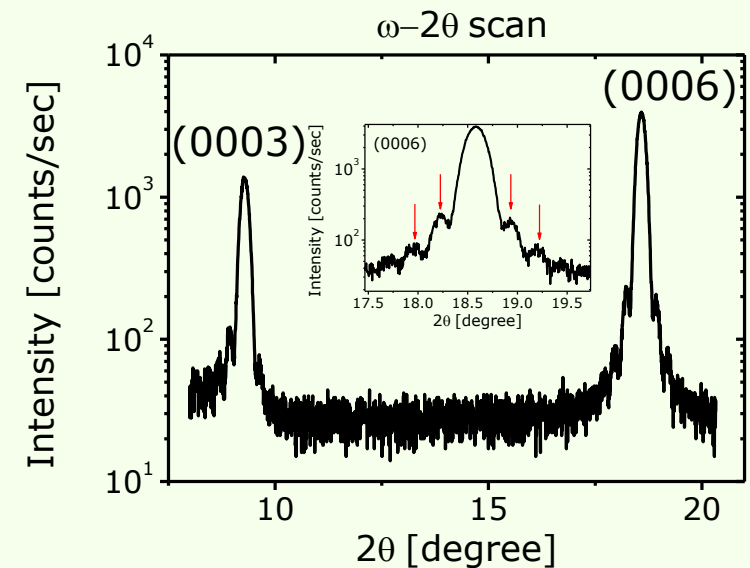
[1,-1,0,0] direction of c-cut sapphire

AFM image shows single quintuple layer steps



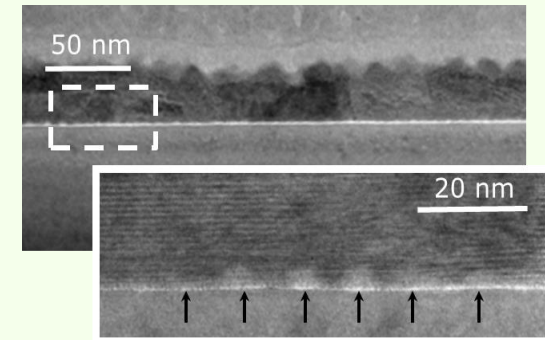
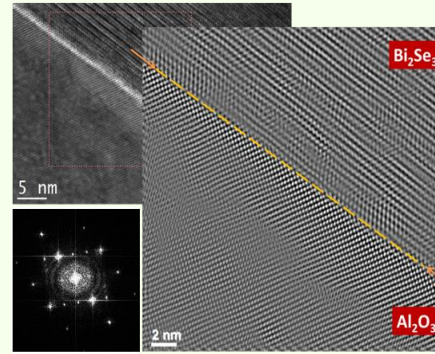
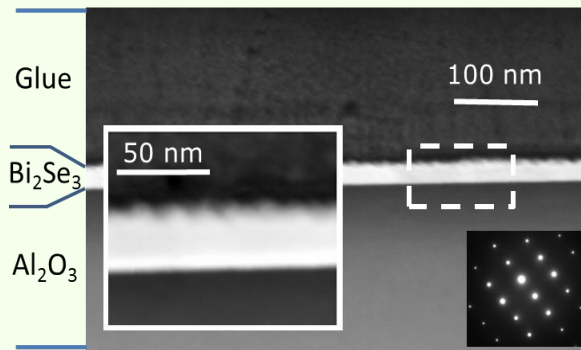
Layer by layer growth
Insitu capping

Heterostruc. possible
TI-I-Metal TI-SC
TI-FMI



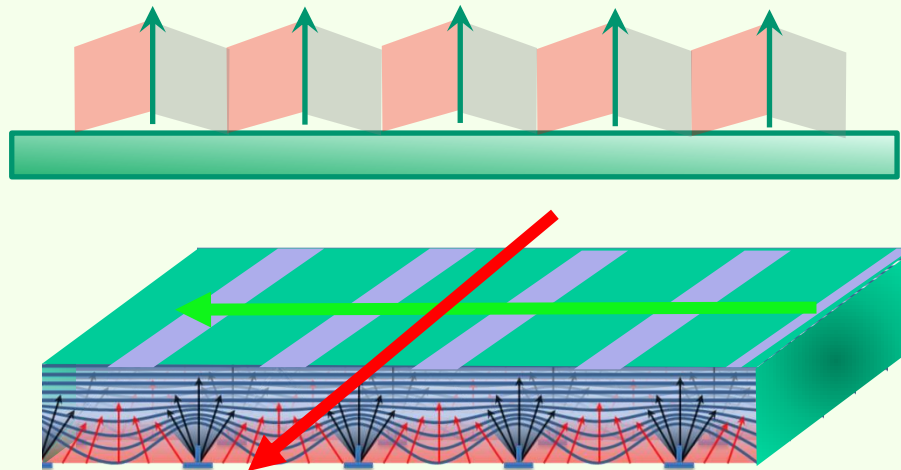
XRD shows single high crystalline quality and terraced surface

- Dislocation mediated structural ordering - Signature of ordering, influence on carrier mobility

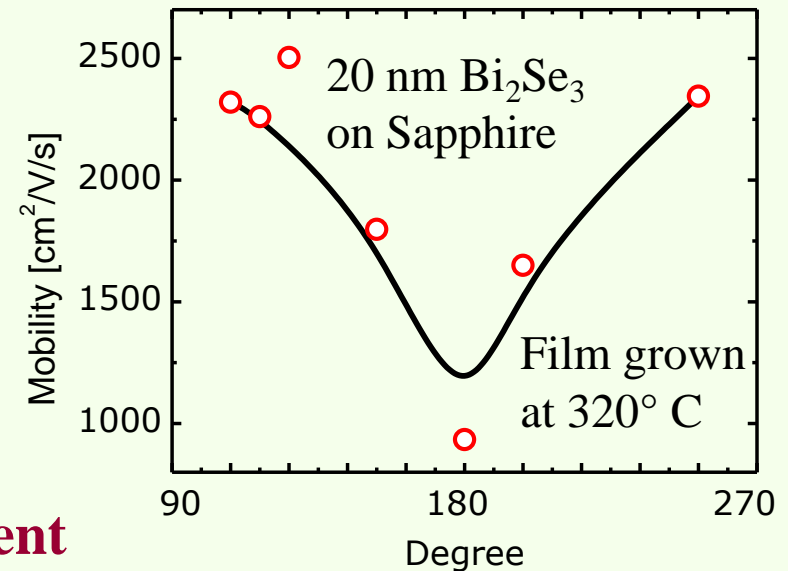


**XTEM, STEM data
very valuable**

Strong structural anisotropy!



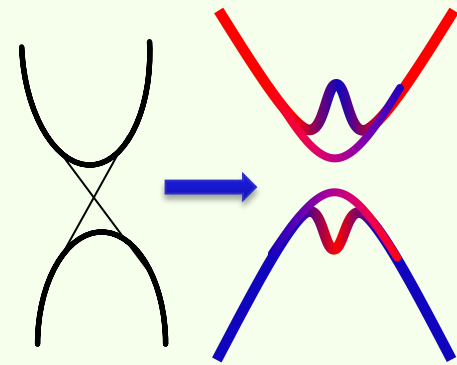
**Mobility depends on the current
direction w.r.t. dislocation lines !**



To observe QAH state

Create a ferromagnetic TI

- Surface Dirac cone opens up a gap



Control film thickness and carrier density to achieve QAHE

Doping TI with magnetic impurities: Cr, Mn, V *etc.*

Reduces surface mobility, spin disorder

Higher T_c requires high doping, lowering $S-O$ - can destroy nontrivial surface states

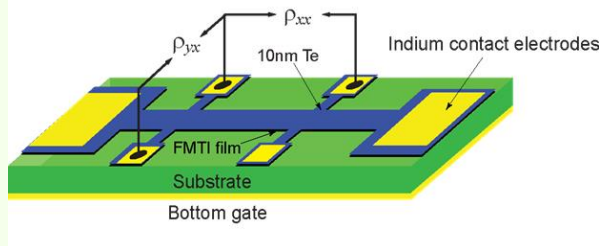
- Proximity driven: through internal exchange field with a magnetic insulator

Independent optimization of elec. and mag. properties

Could be a cleaner approach!

Ferromagnetism in Cr-doped Bi_2Se_3 epitaxial films (2010-11)

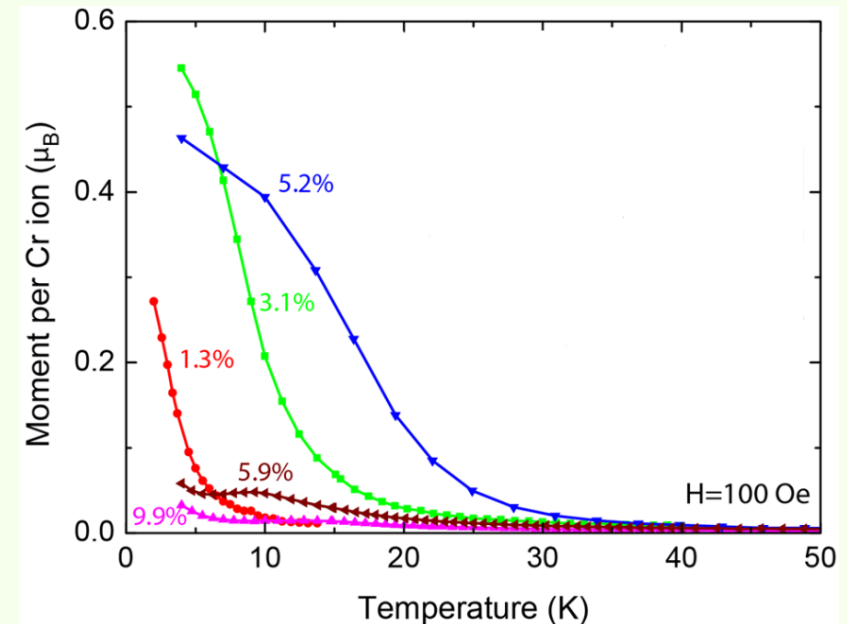
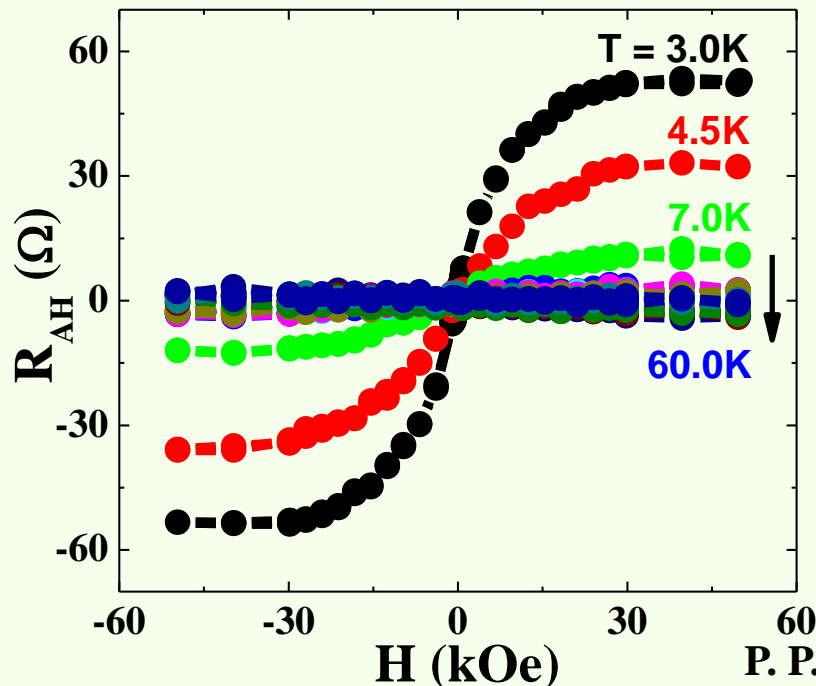
- Lattice parameter c increased ($\sim 0.2\%$)
- Strain increased
- With increasing Cr concentration crystal/film quality decreased



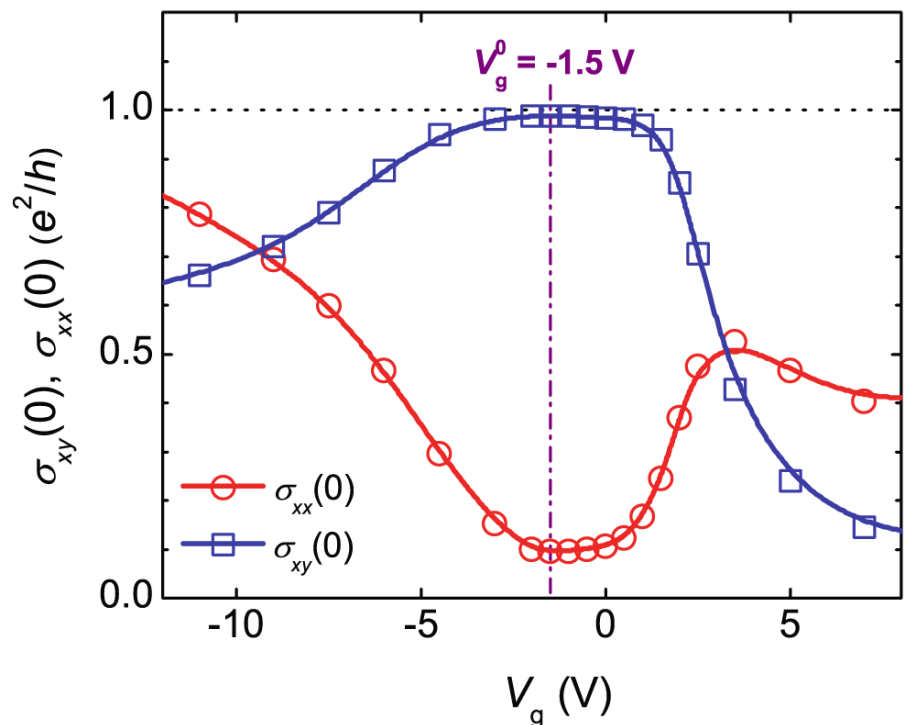
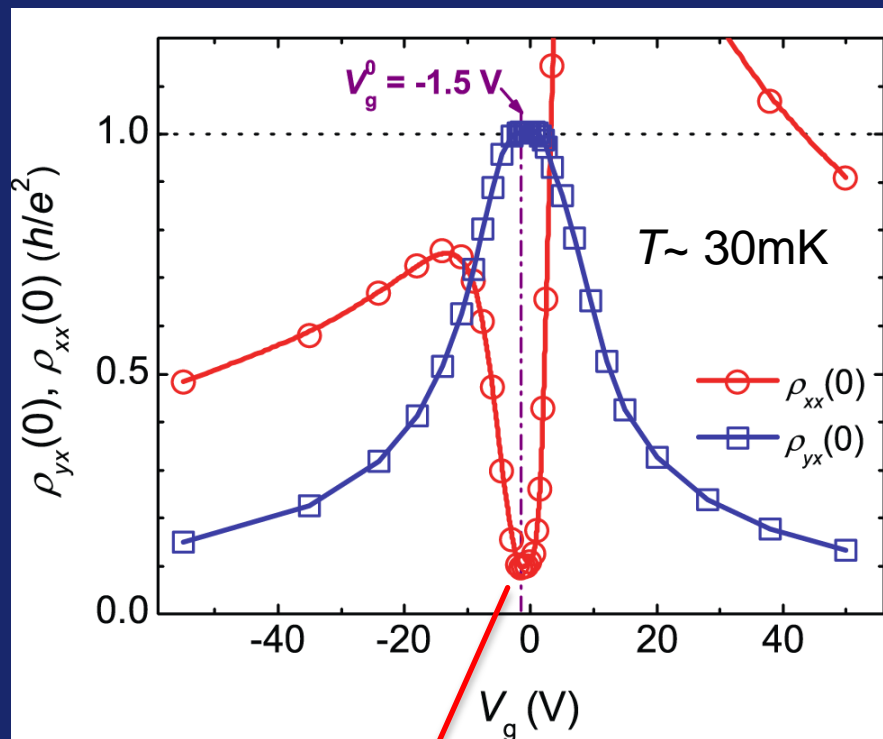
Film quality optimized: 5% Cr

T_c dependence on Cr content

Anomalous Hall effect seen



QAHE in 5QL Cr-doped $(\text{Bi}_{0.1}\text{Sb}_{0.9})_2\text{Te}_3$ film - ferromagnetic TI



Not completely dissipationless

$$\rho_{xx} \sim 2.53 \text{ K}\Omega$$

$$\rho_{yx} = 1 \frac{h}{e^2}$$

$$\sigma_{yx} \sim 0.99 \frac{e^2}{h}$$

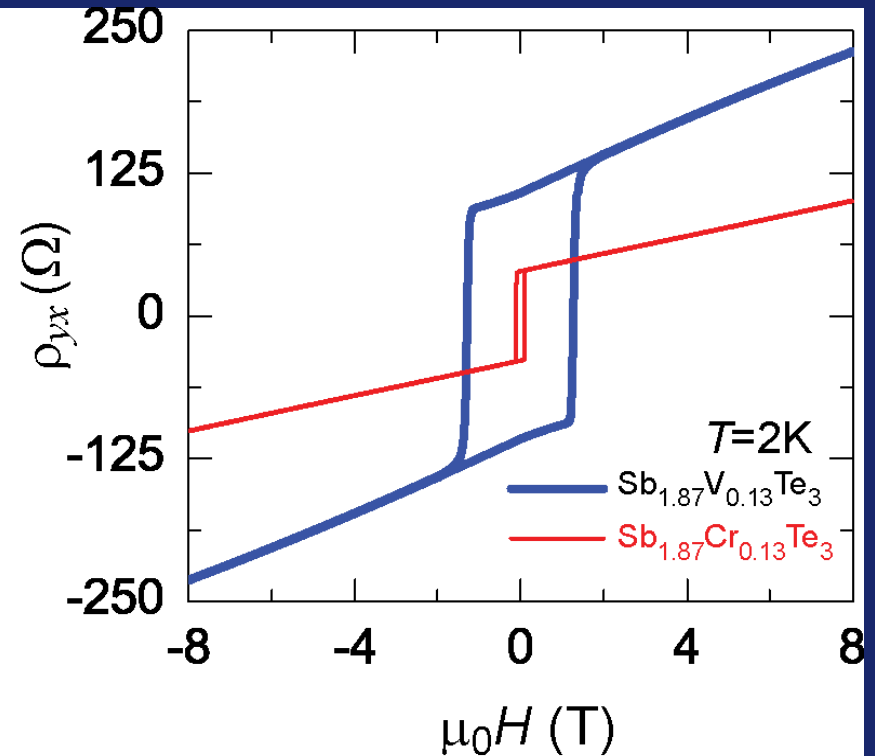
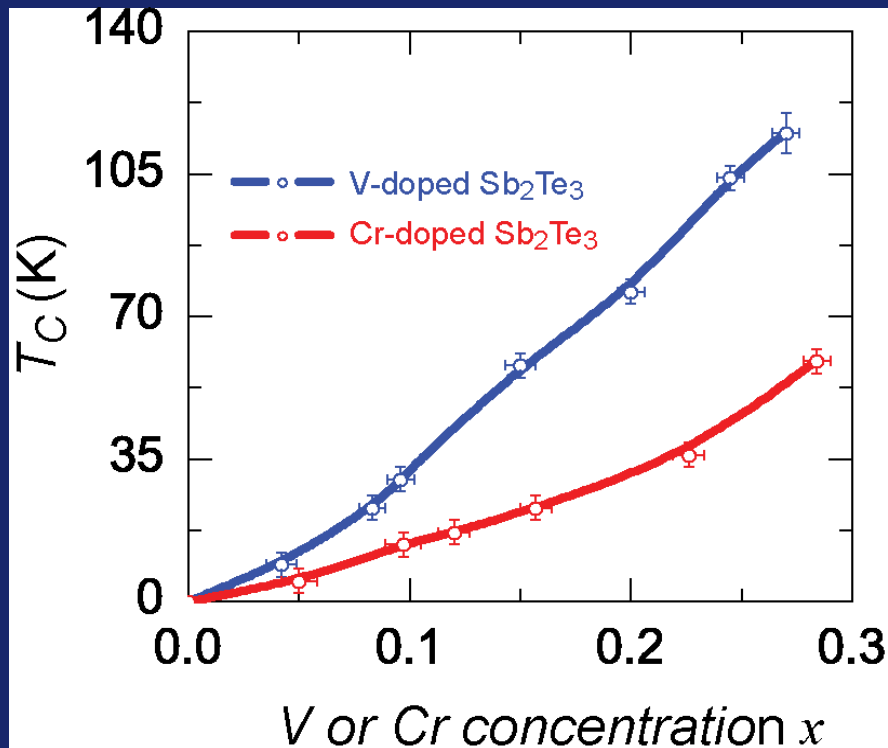
$$\tan \alpha = \frac{\rho_{yx}}{\rho_{xx}} \sim 11$$

Chang et al. Science **340**,167 (2013).

Ferromagnetic properties

Comparing Cr- and V-doped Sb_2Te_3 films

6 QL $\text{Sb}_{2-x}\text{V}_x\text{Te}_3$ and $\text{Sb}_{2-x}\text{Cr}_x\text{Te}_3$

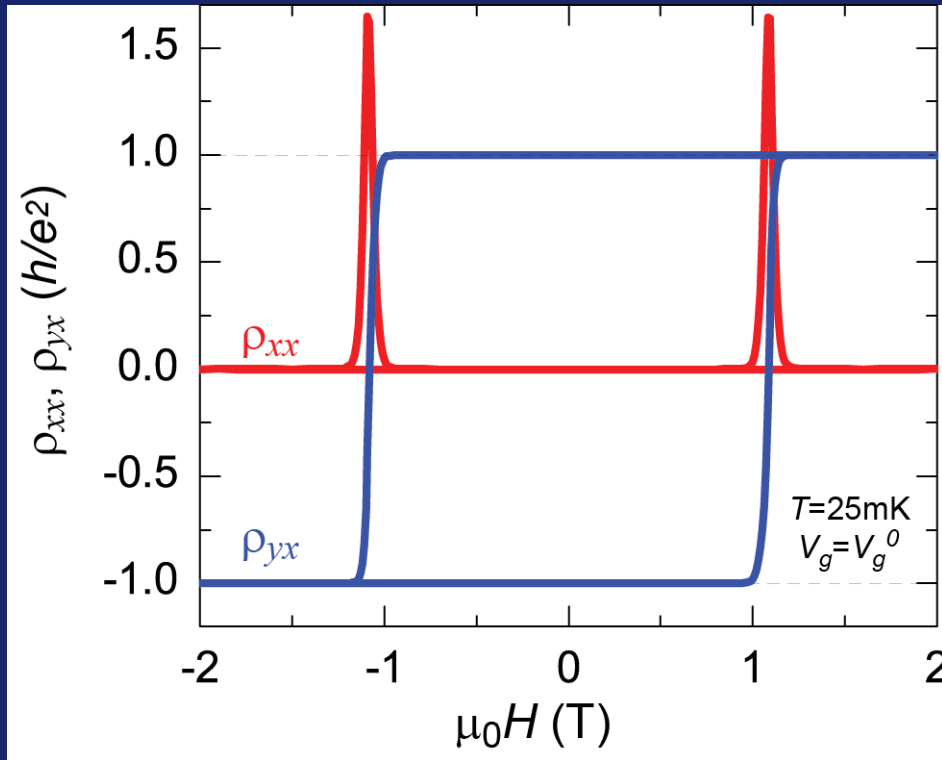


$\text{Sb}_{2-x}\text{V}_x\text{Te}_3$: Higher Curie temp.

Giant coercive field

- Magnetic moment per V ion $1.5 \mu_B$ for $x = 0.13$
 - valence state of V is a mixture of 3+ and 4+ (or/and 5+)
 - reduces p -type carriers by n - doping

High-precision QAHE

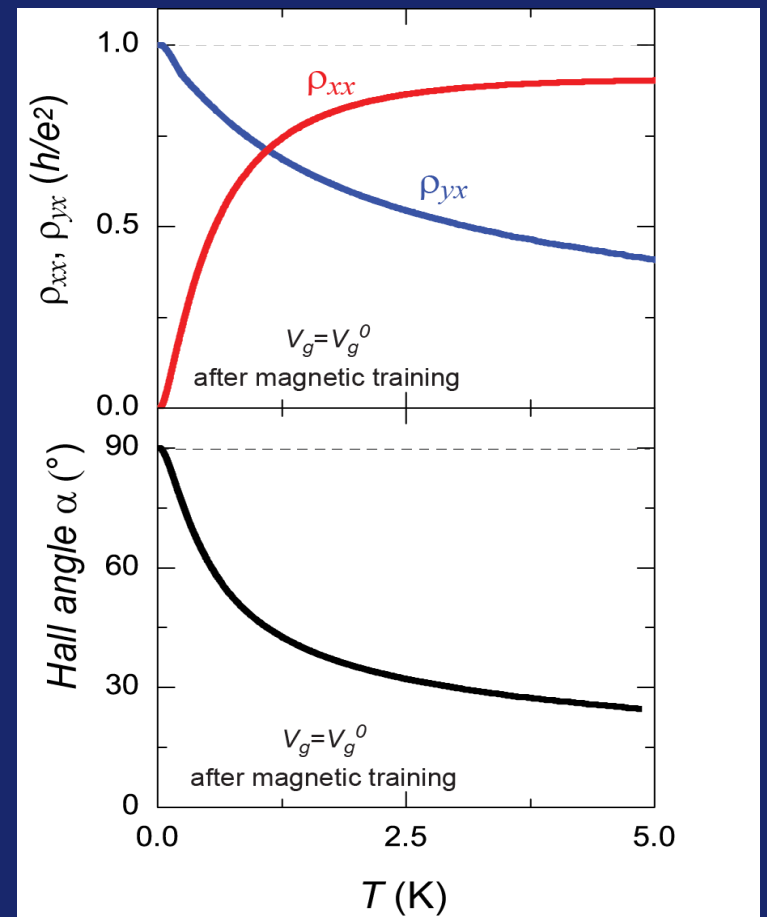


At $H = 0$

$$\rho_{yx} = 1.00019 \pm 0.00069 \text{ } h/e^2$$

$$\rho_{xx} = 0.00013 \pm 0.00007 \text{ } h/e^2 \text{ } (\sim 3.35 \pm 1.76 \Omega)$$

$$\rho_{yx}(0)/\rho_{xx}(0) \sim 7700 \quad \alpha \sim 89.993 \pm 0.004^\circ$$



$$\rho_{yx} = \frac{h}{e^2} = 25.812 \text{ K}\Omega$$

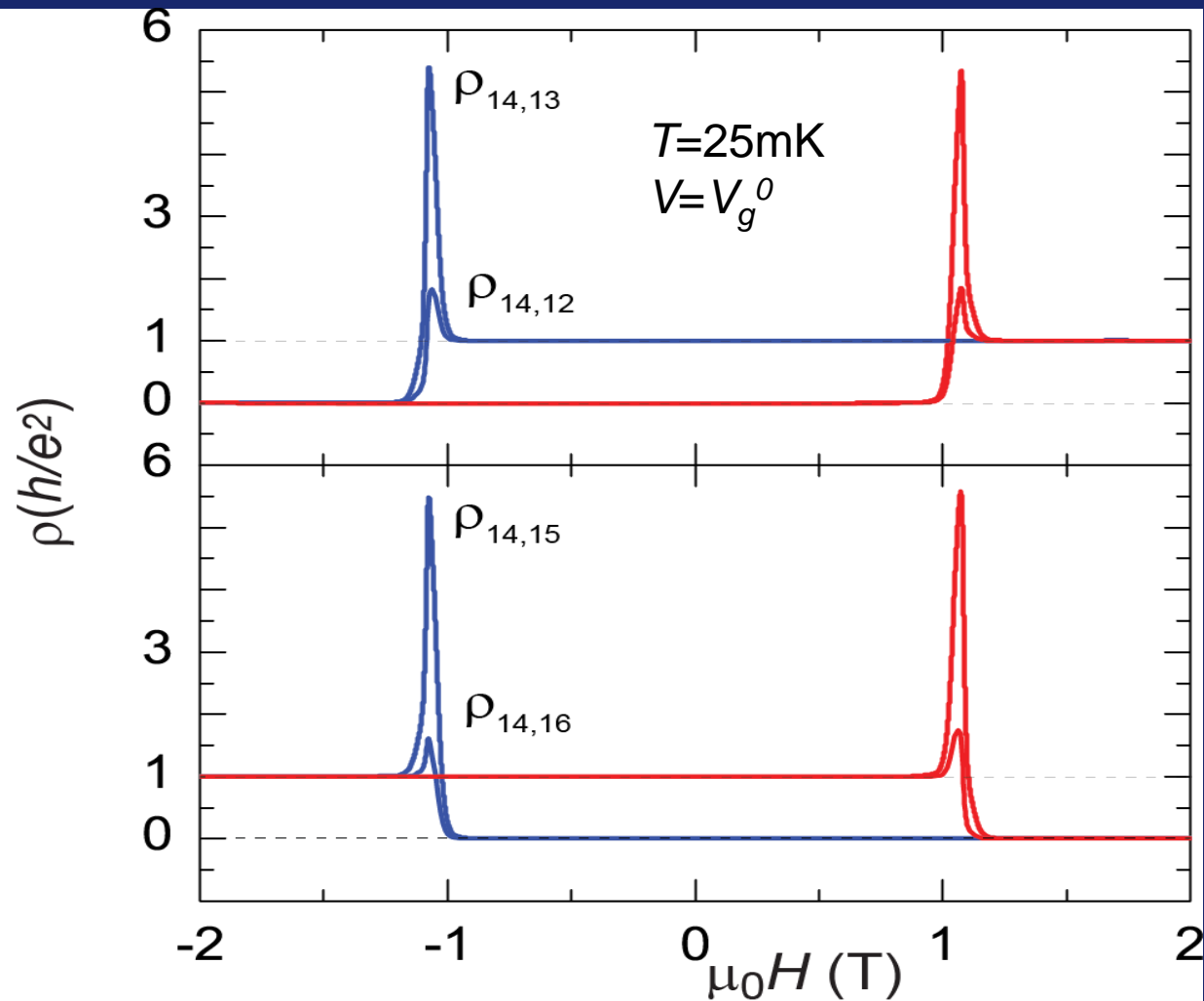
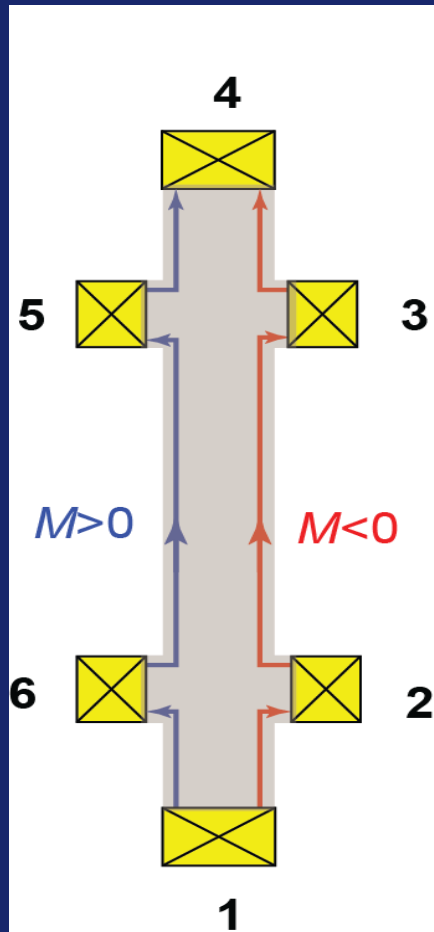
$$\rho_{xx} = 0$$

QAHE state up to $\sim 5\text{K}$

Nat. mater. 14, 473(2015)

Chiral edge dissipationless transport

4QL $V_{0.11}(\text{Bi}_{0.32}\text{Sb}_{0.68})_{1.89}\text{Te}_3$



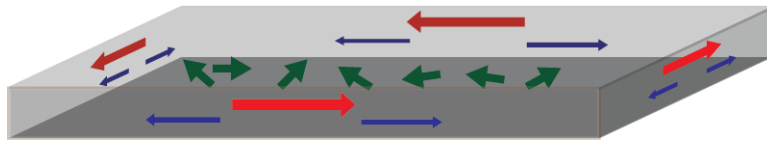
Asymmetric loops

PRL (2015)

Mirror symmetric behavior of $\rho_{14,16}$ ($\rho_{14,15}$) and $\rho_{14,12}$ ($\rho_{14,13}$)

The chirality of edge transport in QAH state

Complex - several conducting channels

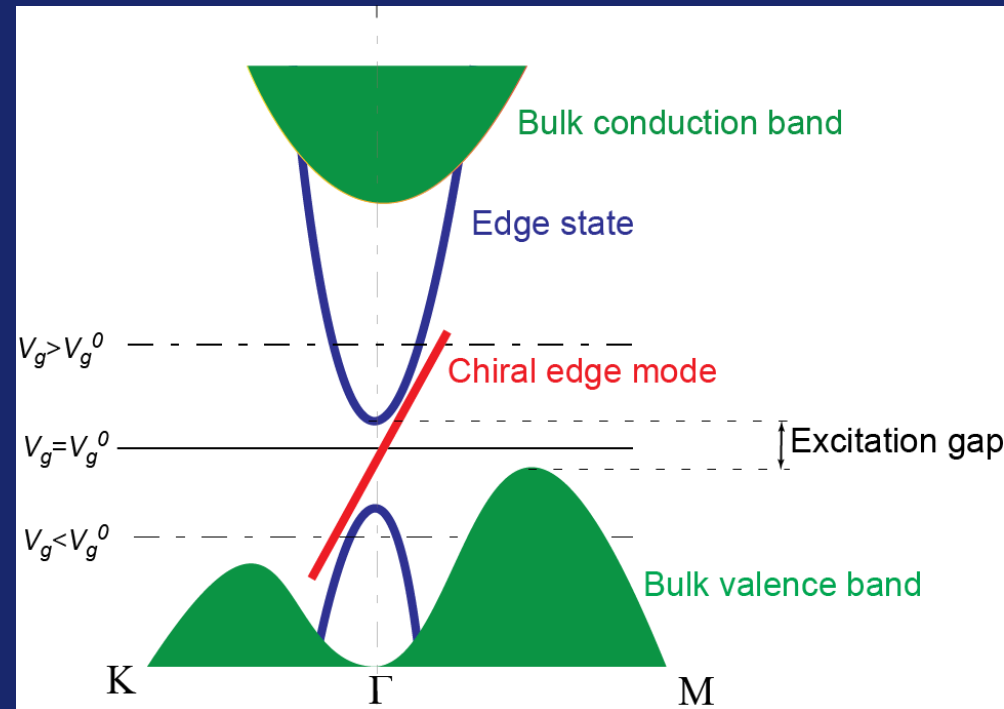


Red: Chiral edge mode

Blue: Nonchiral edge mode

Green: Bulk channels

Non-uniform exchange gap



Activation energy $\sim 100 \mu\text{eV}$

$$V_g = V_g^0$$

Chiral edge channel (red) (dissipationless)

$$V_g > V_g^0$$

Nonchiral edge channel (blue) (dissipative) + Chiral edge channel (red) (dissipationless)

$$V_g < V_g^0$$

Bulk channels (green) (dissipative) + Chiral edge channels (red) (dissipationless)

Creating ferromagnetic state by proximity coupled interfacial exchange field



Very effective with superconductors

PRL 1986, 1988, 1991, 1993, 2007, 2008, 2009;

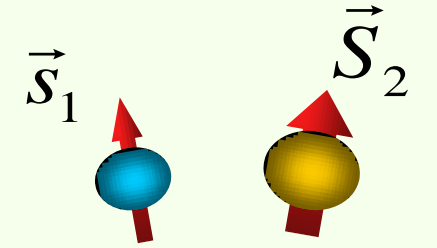
Nat. Commun. 2014; PCCP 2015

Nature 2016, Nature Matl. 2016

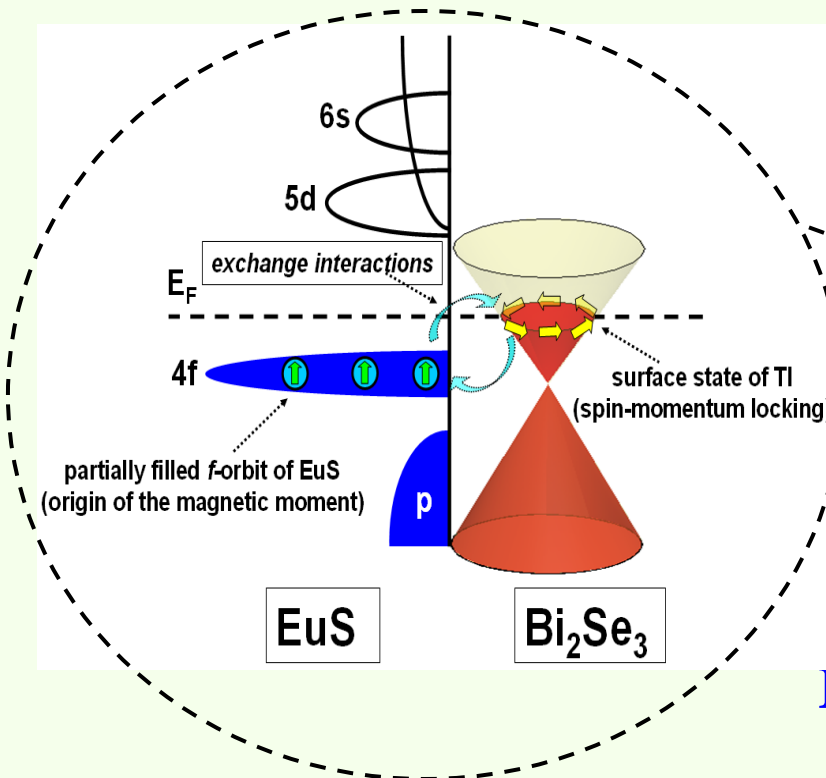
Internal exchange field: Proximity induced surface ferromagnetism in TI

**Large internal exch. field at FI/TI interface
> many Tesla**

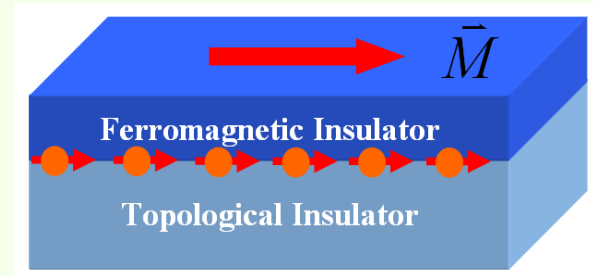
- **Ferromagnetic Insulators (FI): EuS, Eu²⁺ ion: 7 μ_B**
- **4f-5d energy gap: 1.64 eV (EuS) or 1.12 eV (EuO)**
- **Exchange interaction energy J ~ 15 meV**


$$H_{ex} = -J \vec{S}_1 \cdot \vec{S}_2$$

de Gennes 1966
Sarma 1963



FI/TI (bi-layer)

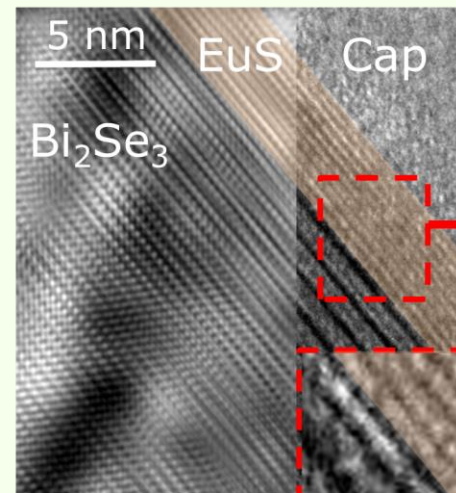


**Split surface bands
(short range interaction)**

Ideal for TI as interaction is at the surface

Introducing ferromagnetism on TI surface via interfacial exchange fields

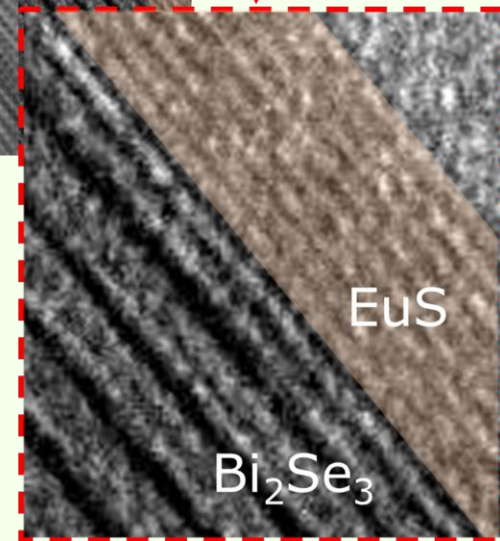
Ferromagnetic Insulator / TI



Epitaxial bilayer of Bi₂Se₃ / EuS
XTEM studies

Independent optimization of elec. and mag.
properties

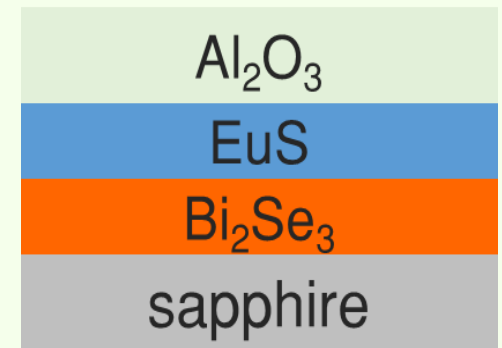
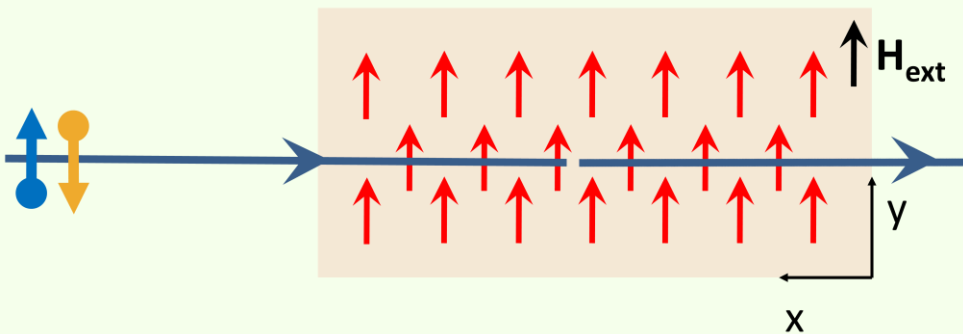
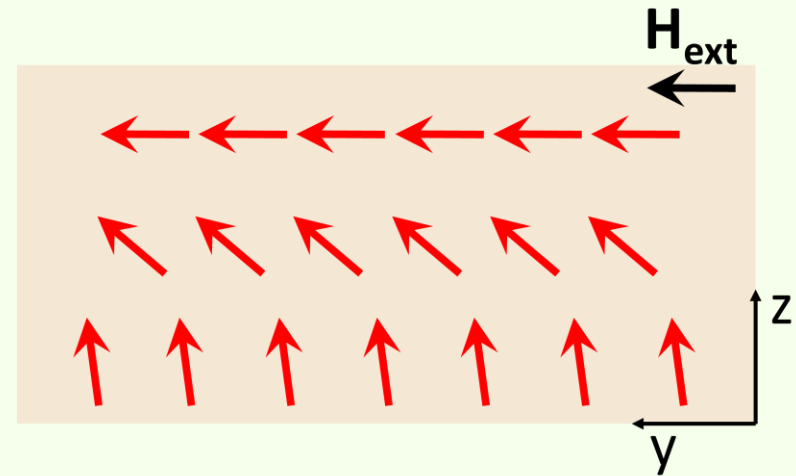
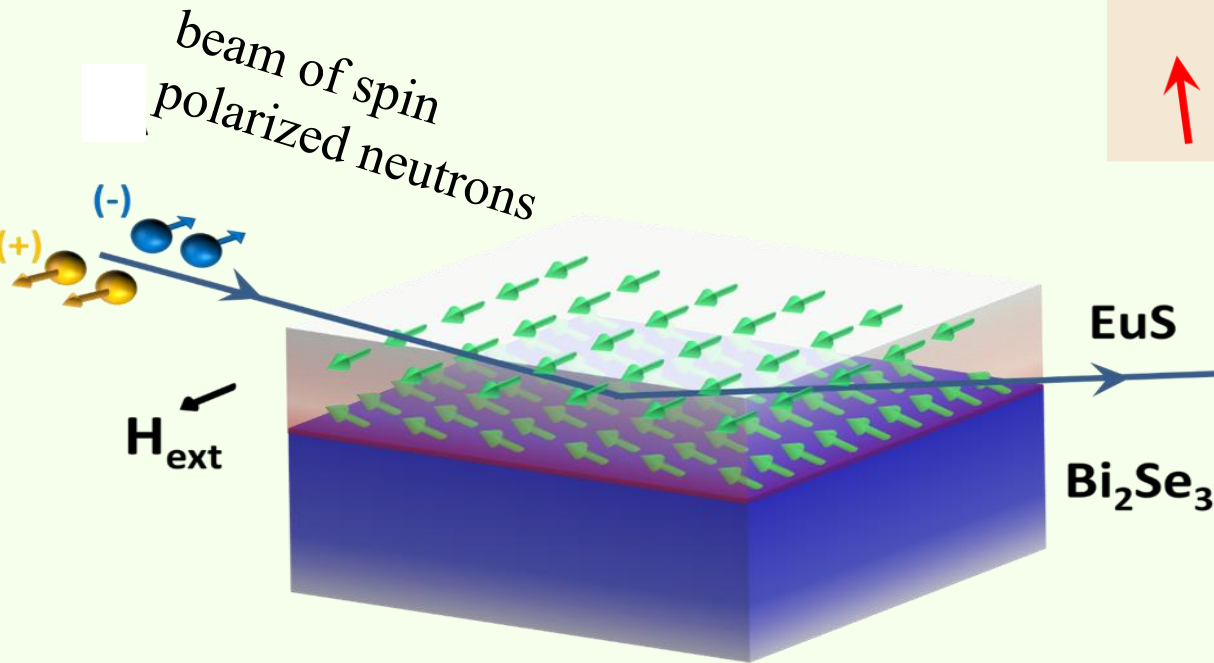
Could be a cleaner approach!



Polarized Neutron Reflectometer

Oakridge Nat. Lab (Valeria Lauter)

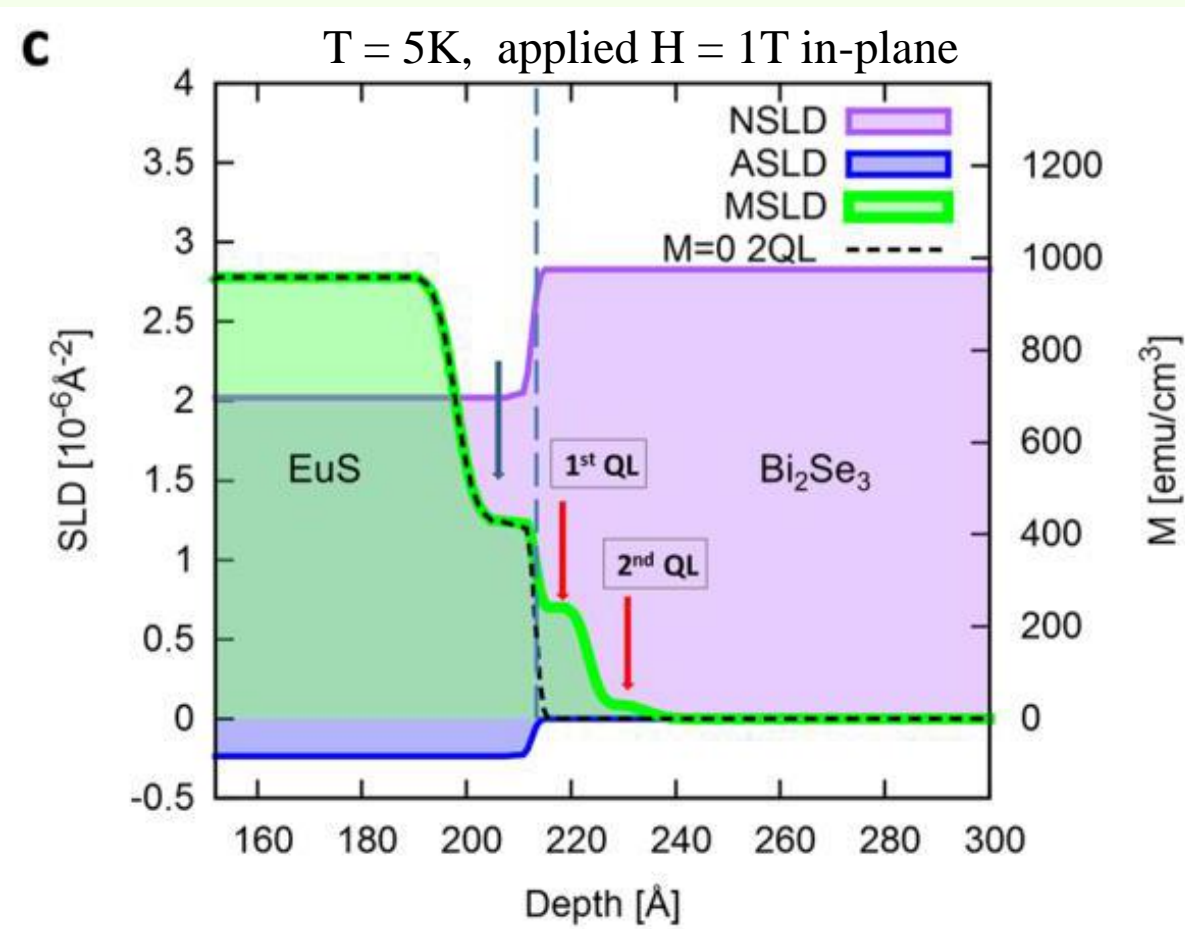
Probing the interface spins



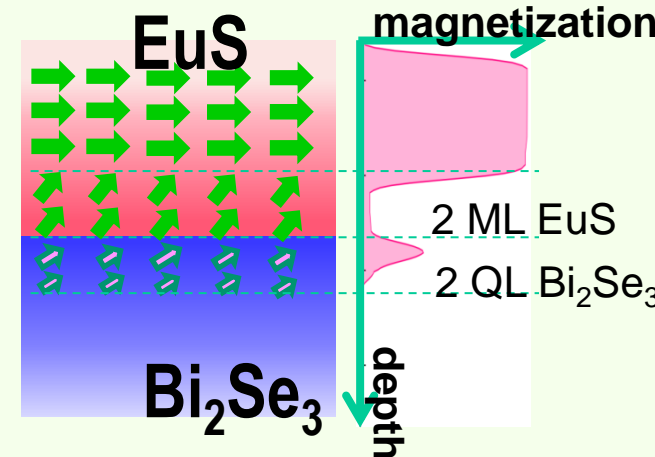
Interface layer ferromagnetism in TI

by Spin-Polarized Neutron Reflectivity

PNR signal showing induced M in EuS/Bi₂Se₃



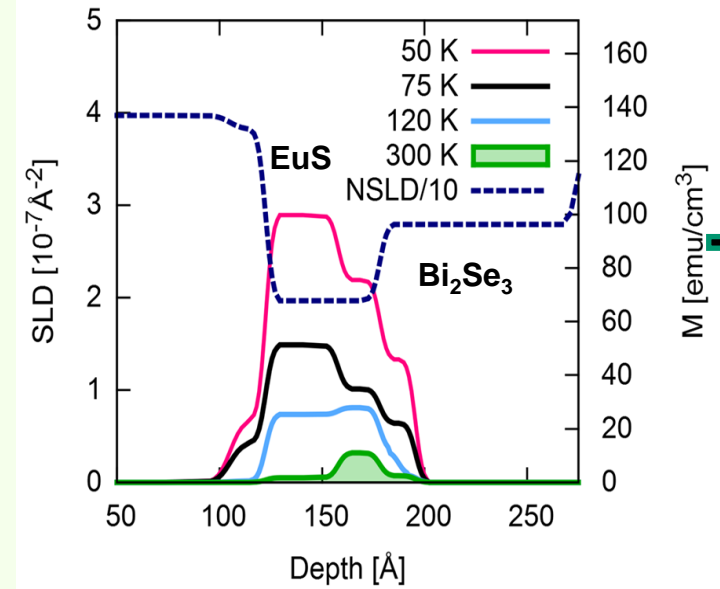
Depth profile of the interface magnetization



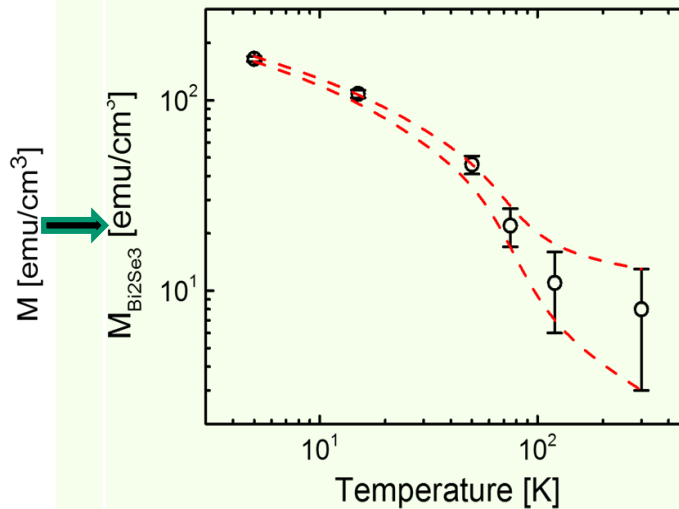
Polarized Neutron Reflectometer: $\text{Bi}_2\text{Se}_3/\text{EuS}$ (10 QL / 5 nm)

Depth profiles at different temp.

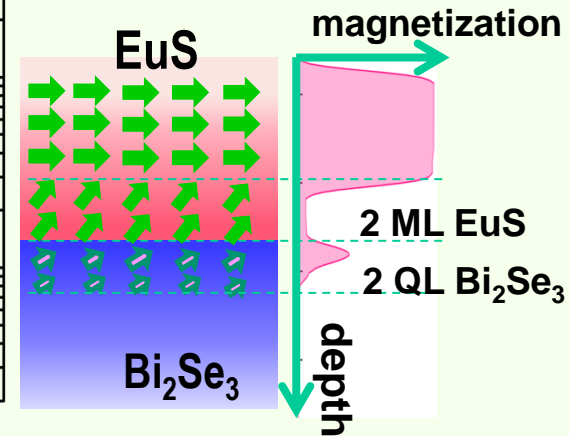
$H_{\text{appl}} = 1\text{T}$ in plane



M vs T of Bi_2Se_3 from PNR data



Depth profile of interface magnetization

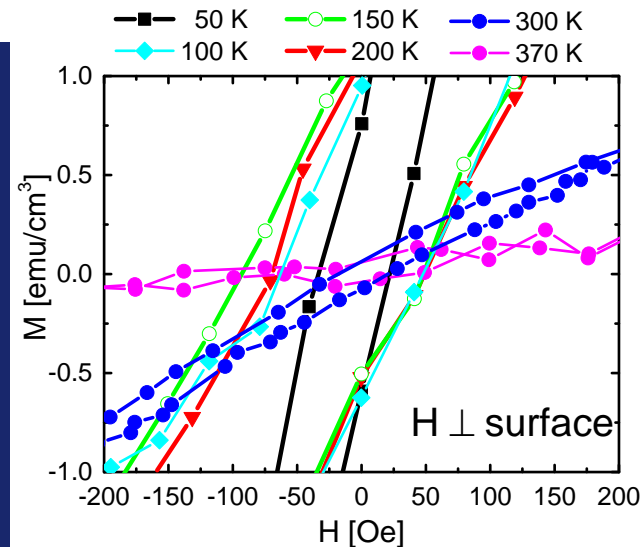
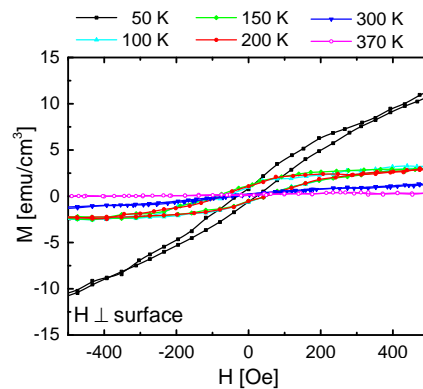
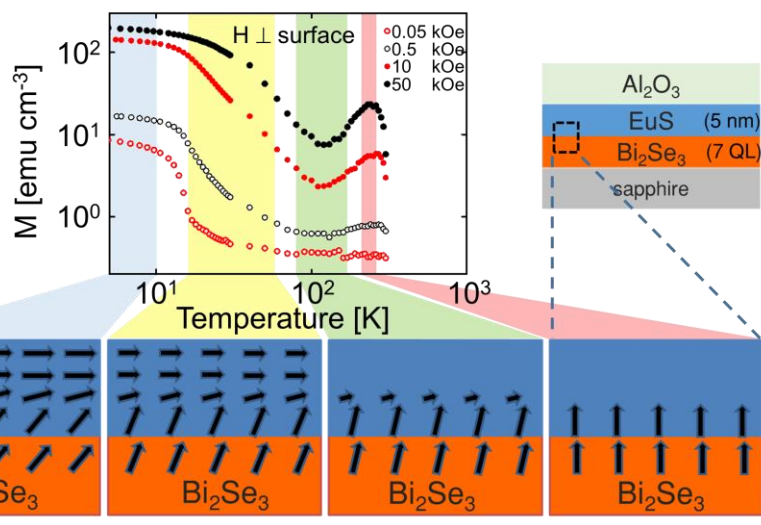


Katmis et al Nature, 2016

- High quality thick EuS layer magnetization confirmed by PNR ($\sim 7m_B$ per Eu^{2+} ion)
- **Canted interface EuS moment (planar $\sim 0.3m_B/\text{Eu}^{2+}$, corresponds to $\sim 85^\circ$ canting angle)**
- **Onset of large induced moment within the top 1-2 QL Bi_2Se_3 , decays with depth**
- No induced interface moment seen in control sample of $\text{EuS}/\text{sapphire}$

SQUID MAGNETOMETRY MEASUREMENT

Magnetization vs. Temp. at various perpendicular applied fields



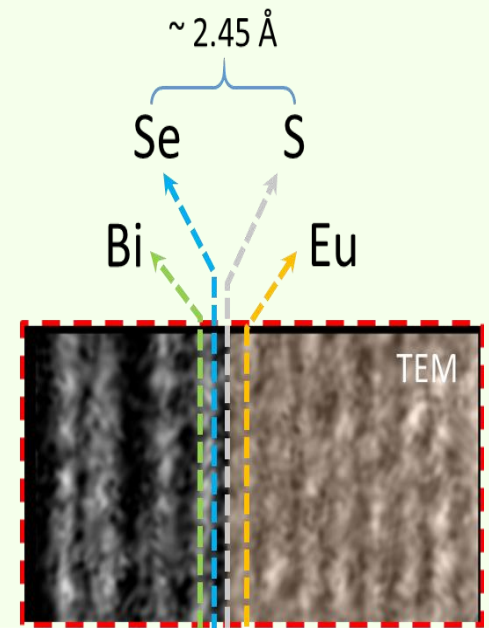
Magnetic hysteresis observed at all T
In agreement with PNR results

At the interface

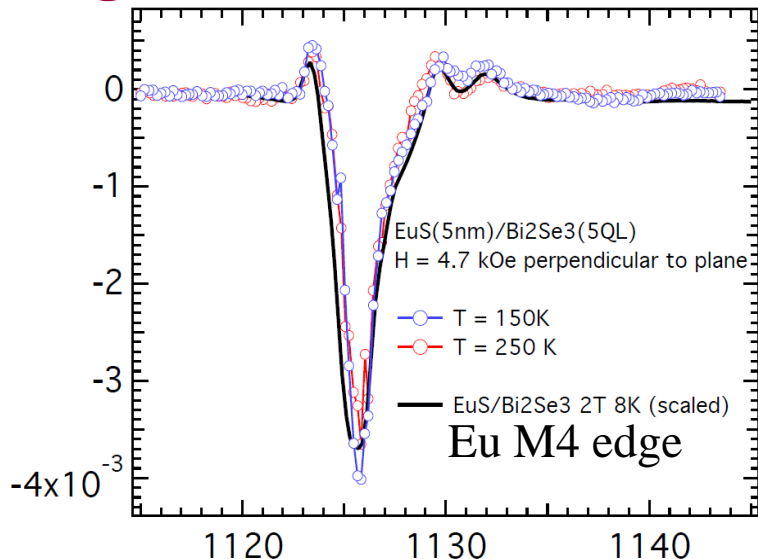
- Large S-O interaction, the spin-momentum locking at Dirac surface state creates strong anisotropy and stabilizes the ferromagnetic state?!

Electron density profile of EuS/Bi₂Se₃ bilayer (glancing angle XRD analysis, synchrotron source used)

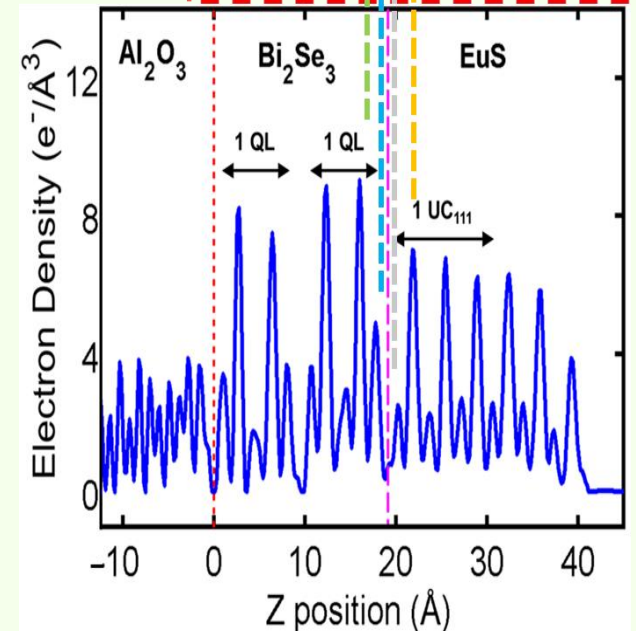
- Bi₂Se₃ is terminated with Se layer at the interface; S layer is in direct stacking with Se layer
- Agrees with TEM analysis
- Sharp interface, no intermixing or inter diffusion
- Weak interfacial bonding van der Waals bond ($\sim 2.45 \text{ \AA}$)



Magnetic circular dichroism vs. Temp.



$\sim 0.05 - 0.1 \text{ uB}$
per atom at 250K

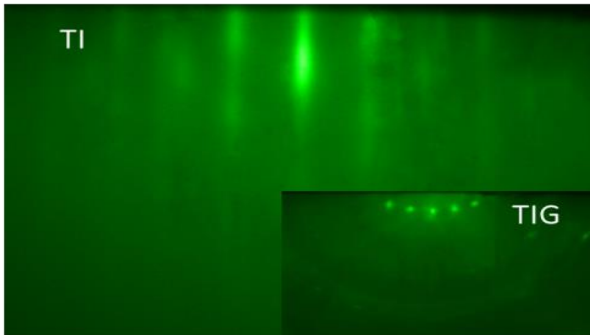


EuS/Bi₂Se₃ interface layer shows ferromag. behavior at 250K $\gg T_C$ (EuS)

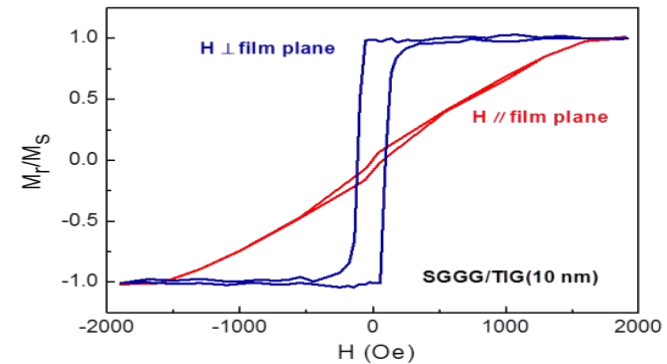
.... Going to higher temperature

Tl/TIG heterostructure (Collab: Jing Shi, UC Riverside)

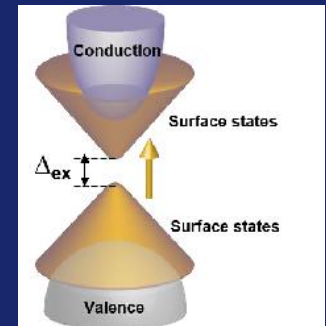
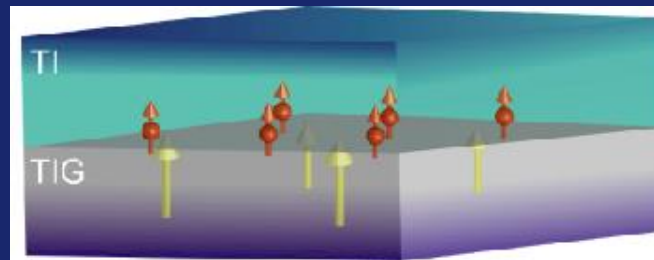
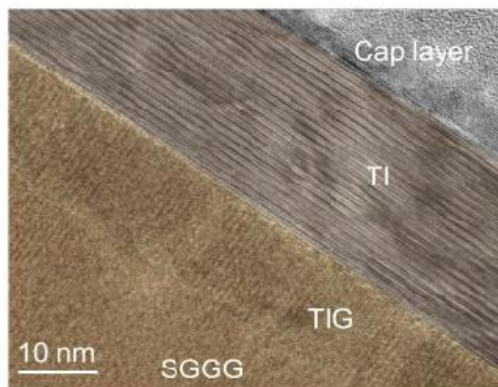
RHEED patterns of Tl/TIG



M-H loop of TIG substrate



TEM image of Tl/TIG



$T_c \sim 560K \rightarrow$ Large exchange-induced gap
Science Advances (Online June 2017)

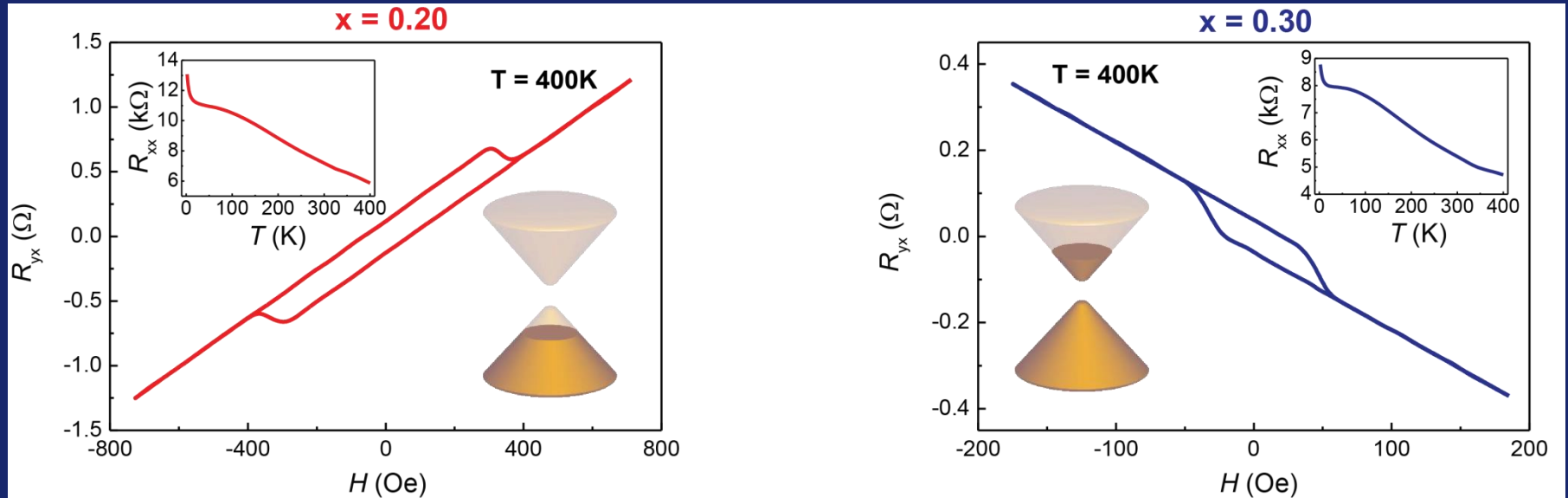
Hall traces of TI/TIG heterostructure

Science Advances (Online June 2017)

5QL $(\text{Bi}_x\text{Sb}_{1-x})_2\text{Te}_3$

p-type

n-type



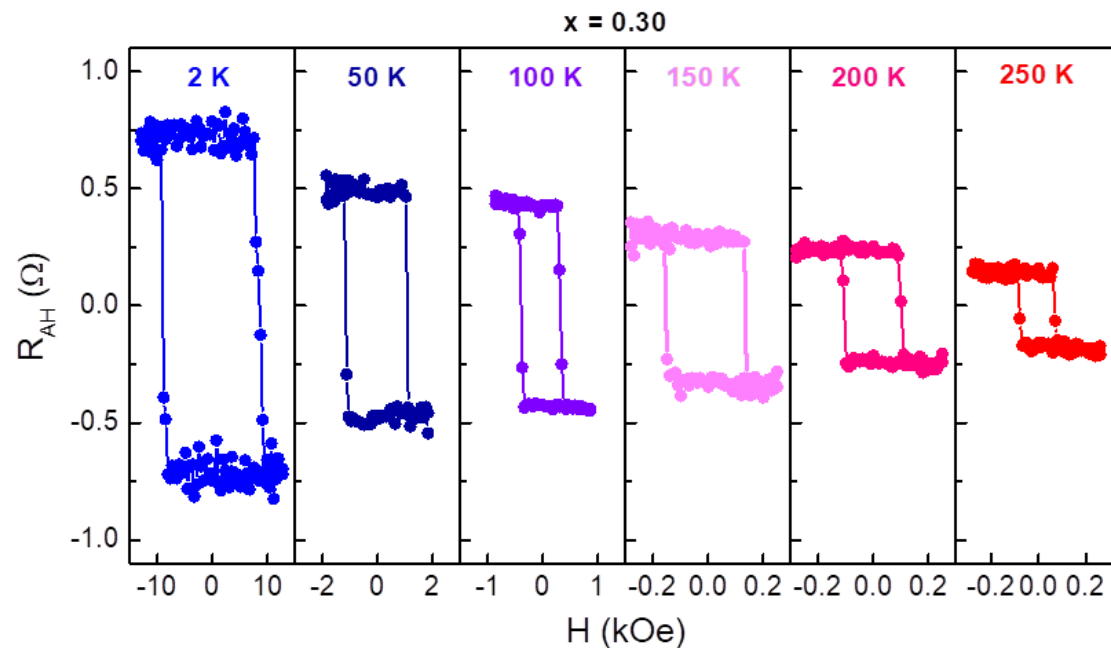
The AHE loops have the same sign in these two samples despite the different carrier types



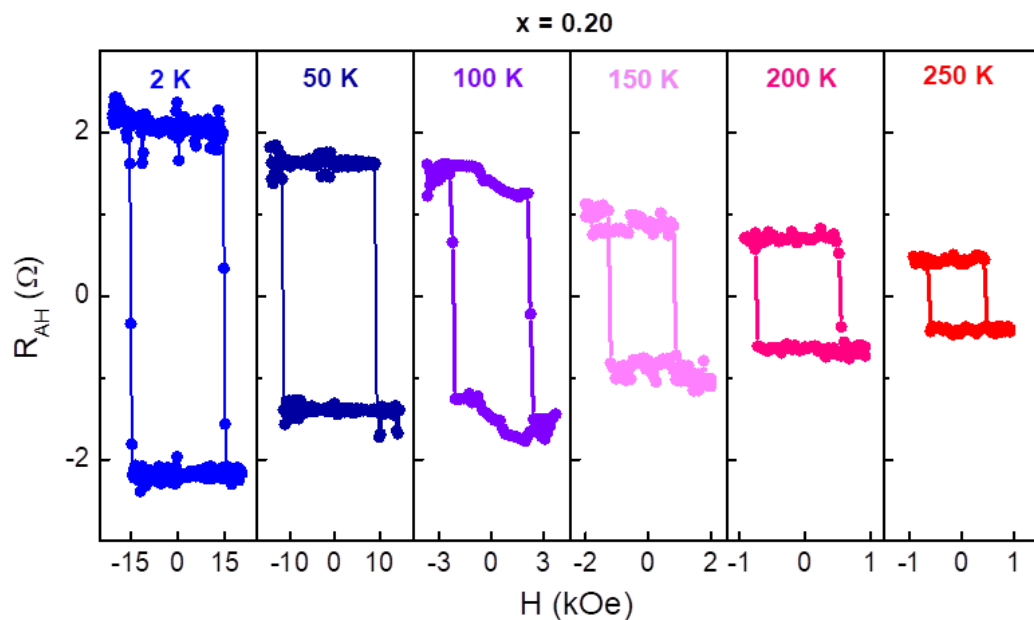
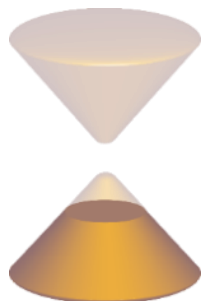
NOT Lorentz force-induced responses arising either from TIG stray fields or contributions from two types of carriers in TI

Anomalous Hall resistance in proximity coupled TIG/ 5QL TI

TIG/(Bi_{0.3}Sb_{0.7})₂Te₃

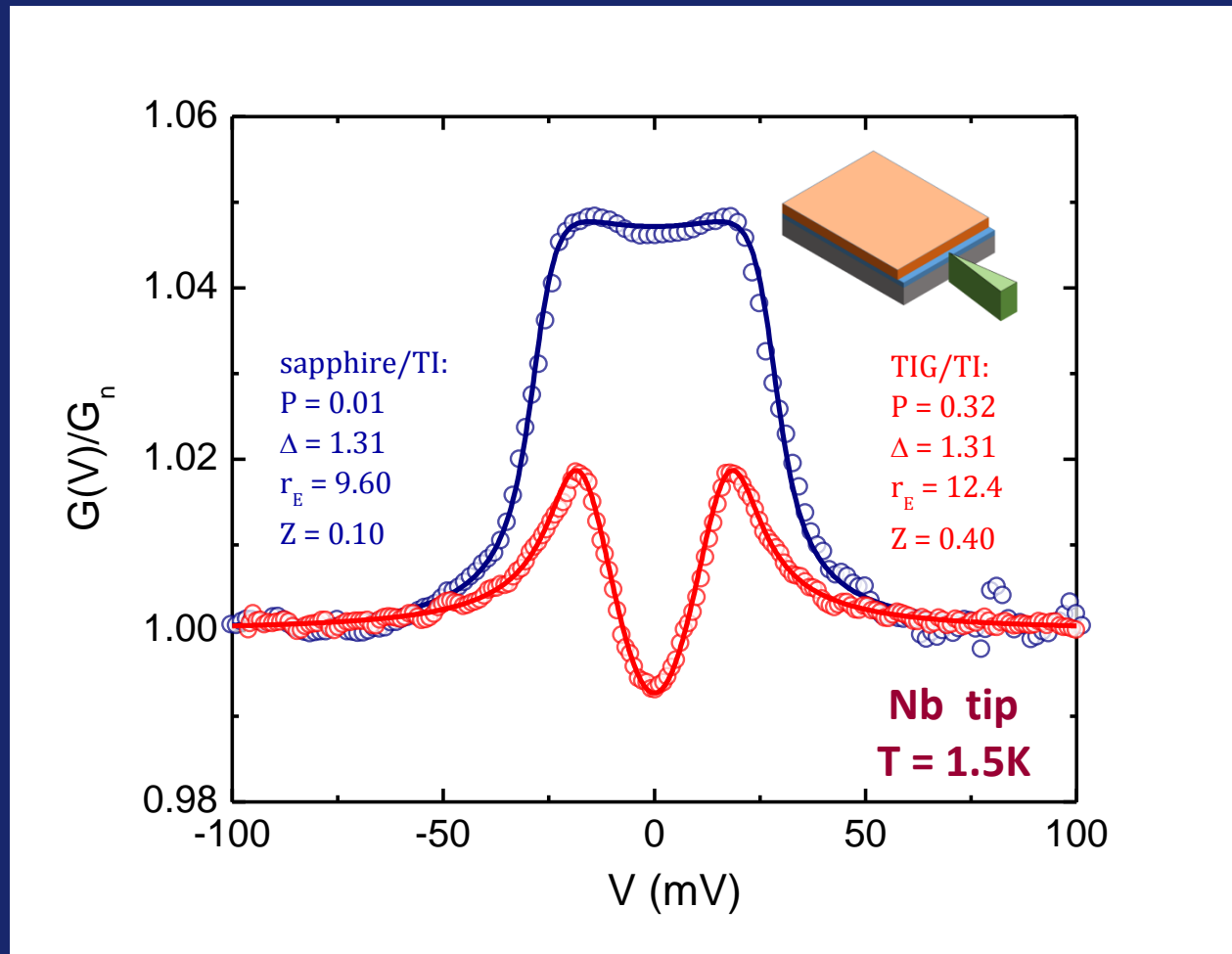


TIG/(Bi_{0.2}Sb_{0.8})₂Te₃



Point-contact Andreev reflection spectroscopy (PCAR) on TI/TIG heterostructure

Science Advances (Online June 2017)



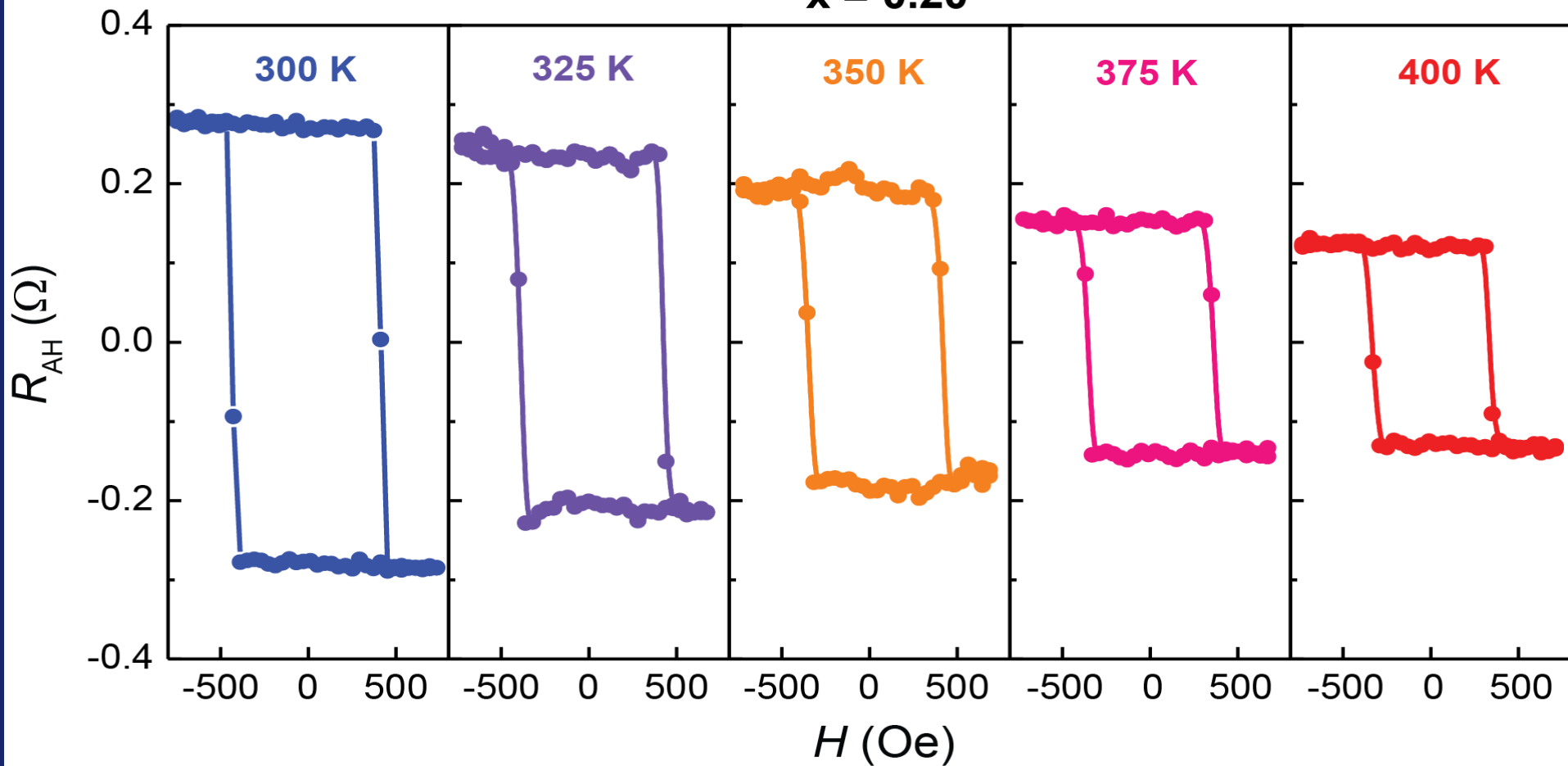
Spin polarization at the interface of TI/TIG heterostructure,
indicates the interfacial induced FM phase in TI

Temperature dependence of AHE response in TI/TIG

5QL $(\text{Bi}_x\text{Sb}_{1-x})_2\text{Te}_3$

p-type

$x = 0.20$



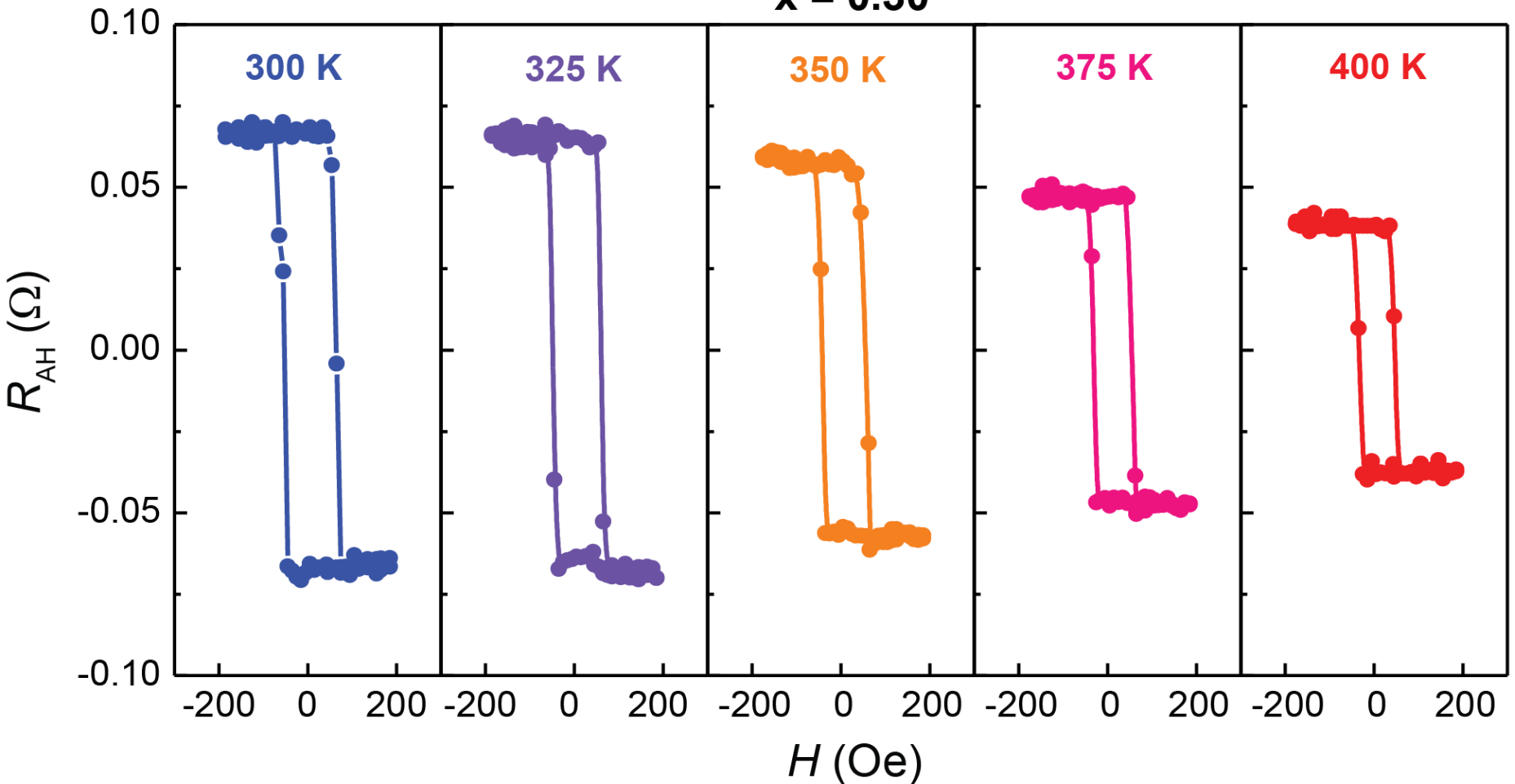
- Perpendicular FM phase exists throughout the entire temperature range up to 400 K
- AHE resistance increase with decreasing the temperature

Temperature dependence of AHE response in TI/TIG

5QL $(\text{Bi}_x\text{Sb}_{1-x})_2\text{Te}_3$

n-type

$x = 0.30$



- Perpendicular FM phase exists throughout the entire temperature range up to 400 K
- AHE resistance increase with decreasing the temperature

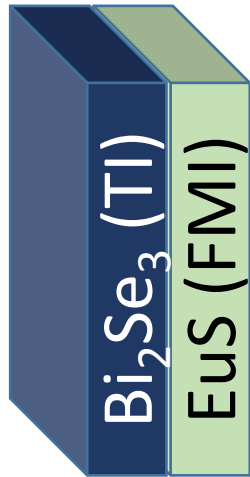
Heterostructure FMI/TI bilayer

– magnetic order by surface states

(Kyungmin Lee, Nandini Trivedi)

Ohio State Univ.

Model:

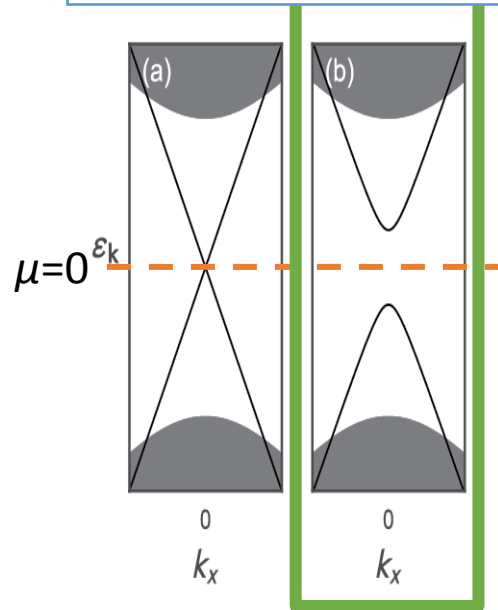


Kondo coupling
between spin and
surface state

$$\mathcal{H}_{\text{Kondo}} = J \sum_{\mathbf{x}} \mathbf{S}(\mathbf{x}) \cdot \mathbf{c}_{\mathbf{x}}^{\dagger} \boldsymbol{\sigma} \mathbf{c}_{\mathbf{x}}$$

generates effective
spin-spin interaction
through RKKY-like
mechanism

Effect of **FMI** on
TI: Gapping of
topological
surface states of
TI

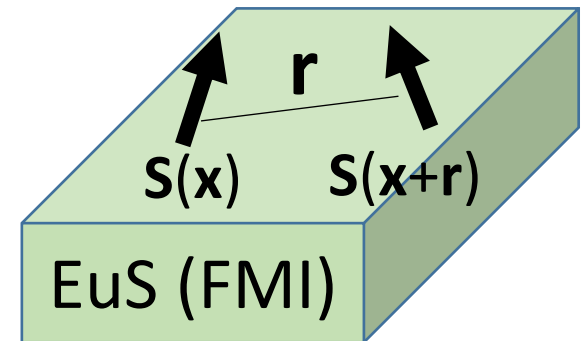


Before
coupling
to FM:
GAPLESS

After
coupling to
FM: **Massive
Dirac**

χ : spin-spin correlation of surface states

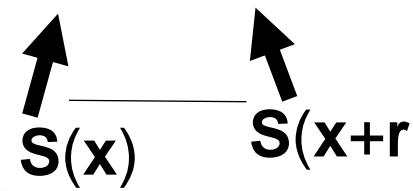
Effect of **TI** on **FMI**:
Additional exchange
and DM coupling
generated between
classical spins



$$\delta F_2[\mathbf{S}] = -J^2 \sum_{\mathbf{x}, \mathbf{y}} S_i(\mathbf{x}) S_j(\mathbf{y}) \chi_{ij}(\mathbf{x} - \mathbf{y})$$

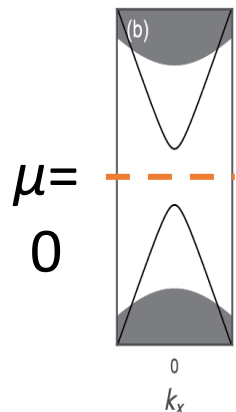
$$\tilde{J}_{ij}(\mathbf{r}) = -J^2 \chi_{ij}(\mathbf{r})$$

Directional dependence of effective spin interaction

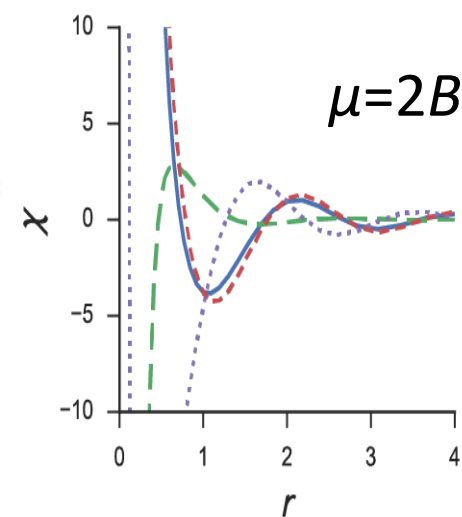
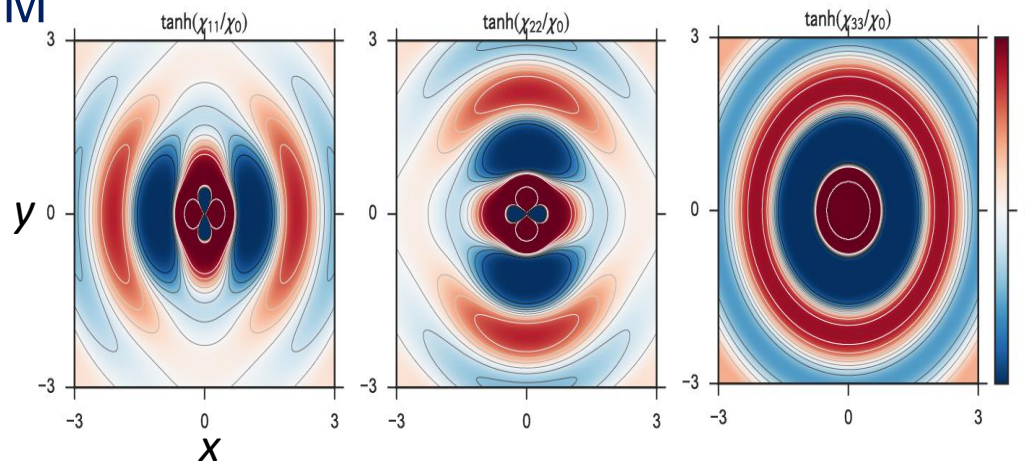
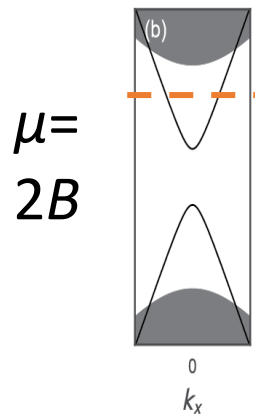
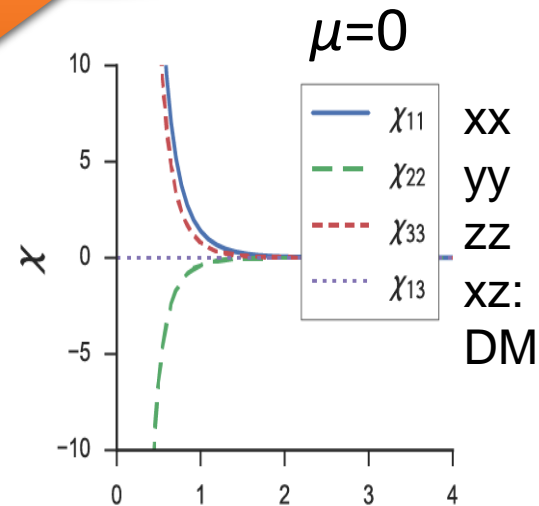
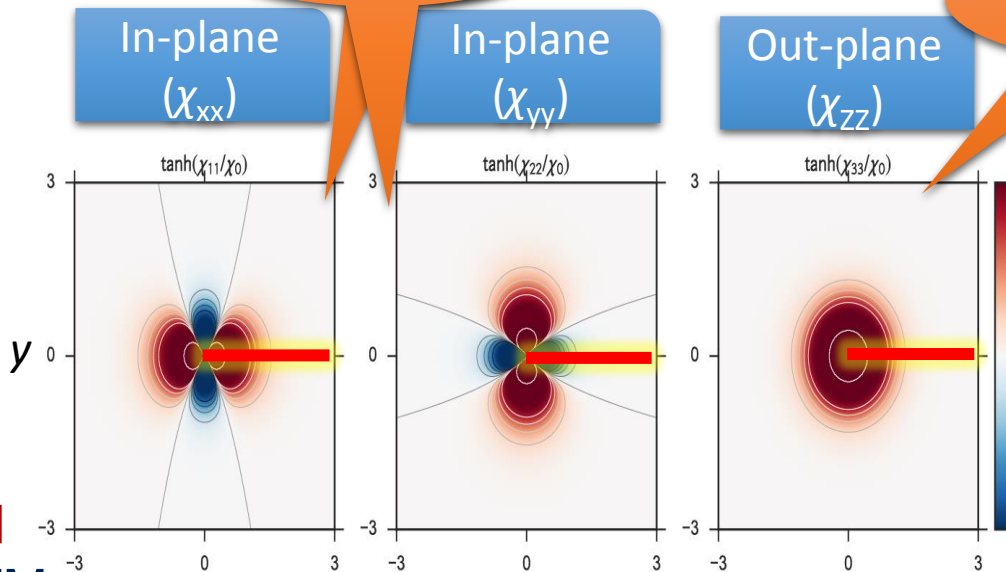


Highly anisotropic

Ferromag.



Red: FM
Blue: AFM

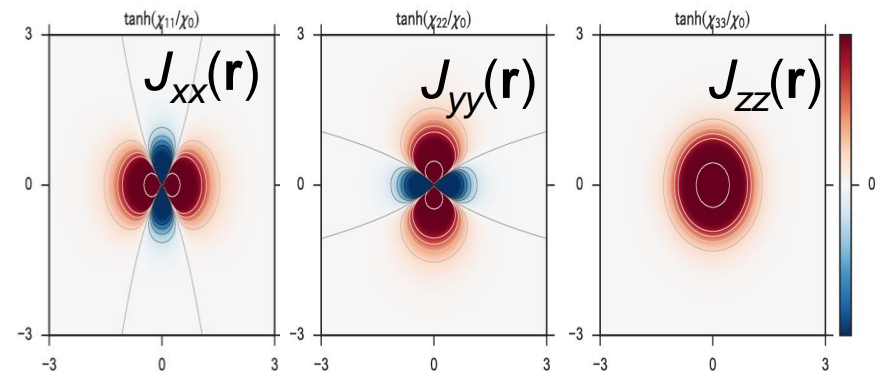


Results for effective spin-spin interaction

- Effective spin interaction in FM mediated by the TI surface states can enhance T_C .
- Spin-momentum locking of surface state \rightarrow highly anisotropic spin interaction (RKKY of spin-degenerate metal is isotropic)
- “Frustrated” in-plane vs. stable ferromagnetic out-of-plane interaction \rightarrow consistent with out-of-plane magnetism at interface – exptl. obser.
- Oscillation in spin interaction as function of distance \rightarrow Possibly sensitive to lattice constant change
- Both Heisenberg and Dzyaloshinskii-Moriya present \rightarrow Possible skyrmion phase(?)

Result summary

	(Anisotropic) Heisenberg	Dzyaloshinskii-Moriya
$\mu=0$	<ul style="list-style-type: none"> - Longitudinal: FM - In-plane transverse: AFM - Out-of-plane: FM 	$J_{DM}=0$
$\mu=2$ B	<ul style="list-style-type: none"> - Friedel oscillation - Same as $\mu=0$ at short distances 	$J_{DM} \neq 0$



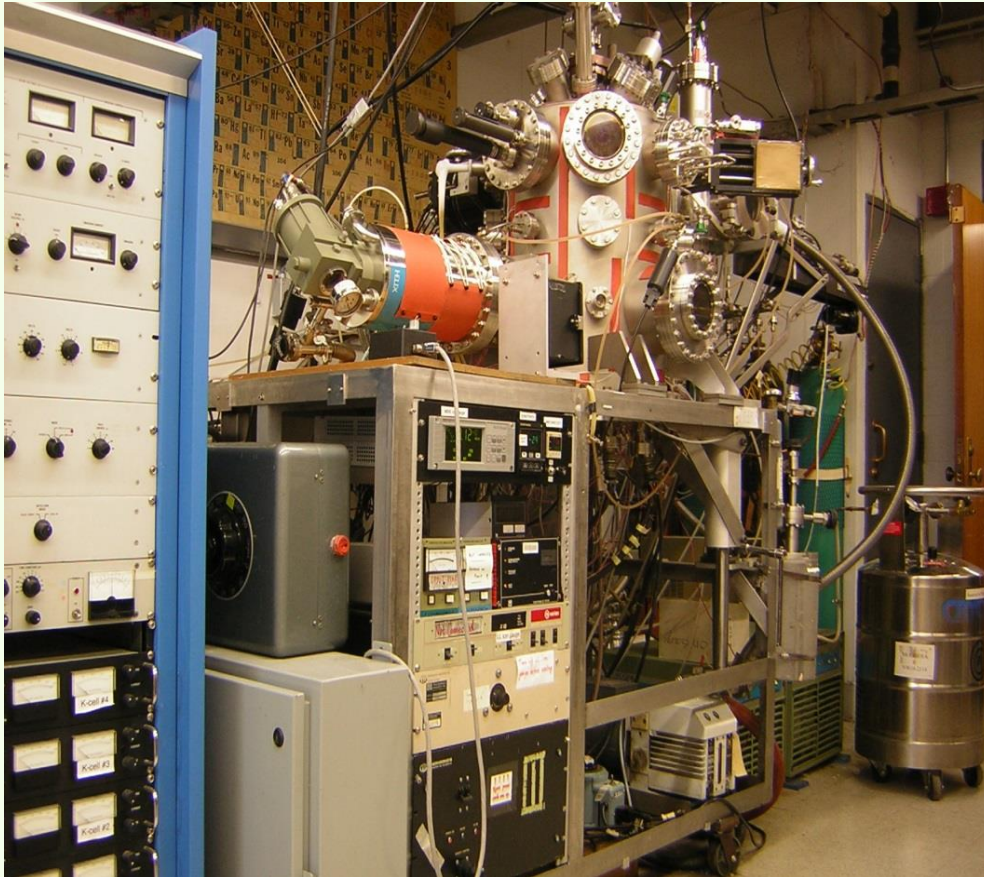
Summary

Investigating TI quantum properties

REQUIRES high quality TI films and heterostructures

- **Ferromagnetic TI - *transition metal doped***
QAH state obtained
Chiral dissipationless edge conduction observed
- **- *TI proximity coupled to ferromag. Insul. - works***
Domain wall chiral conduction observed
Ferromagnetic TI at RT – interface driven
QAH state is yet to be reached

Custom Built Metal MBE



- 10^{-10} Torr, 14 sources, masks, plasma, insitu multilayers, sputtering,..... in-situ RHEED
- heating/cooling (80K to 1200K).....
- angled deposition – nano structure creation

Carl Zeiss Leo Supra 25

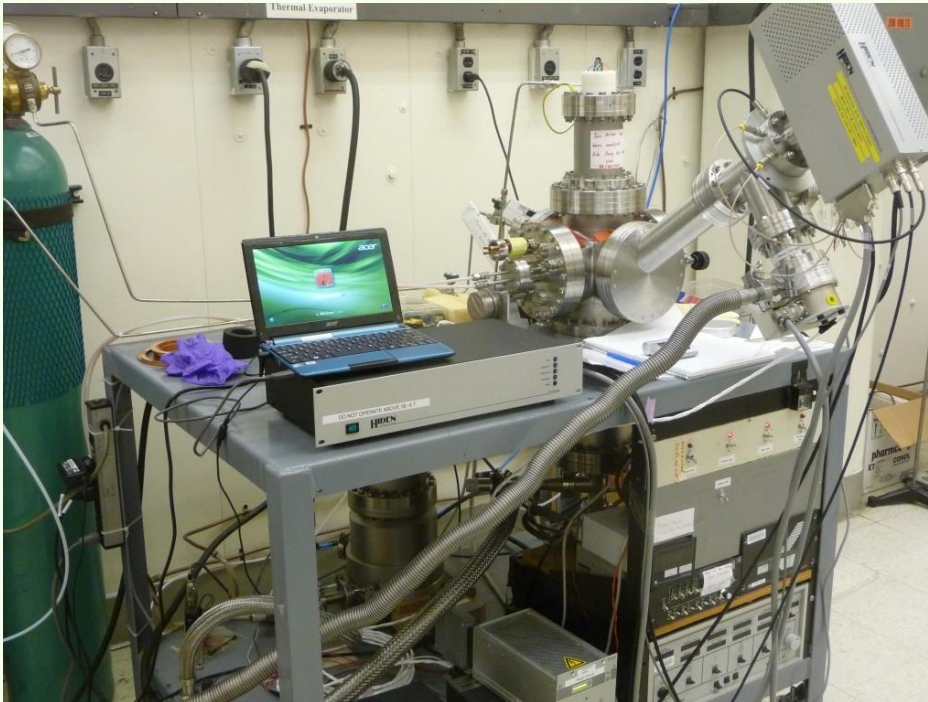


SEM (2nm)

Raith ebeam writing -
nanopatterning (20nm
features)

EDX – elemental analysis

Nano patterning: Ion Miller with end point detector



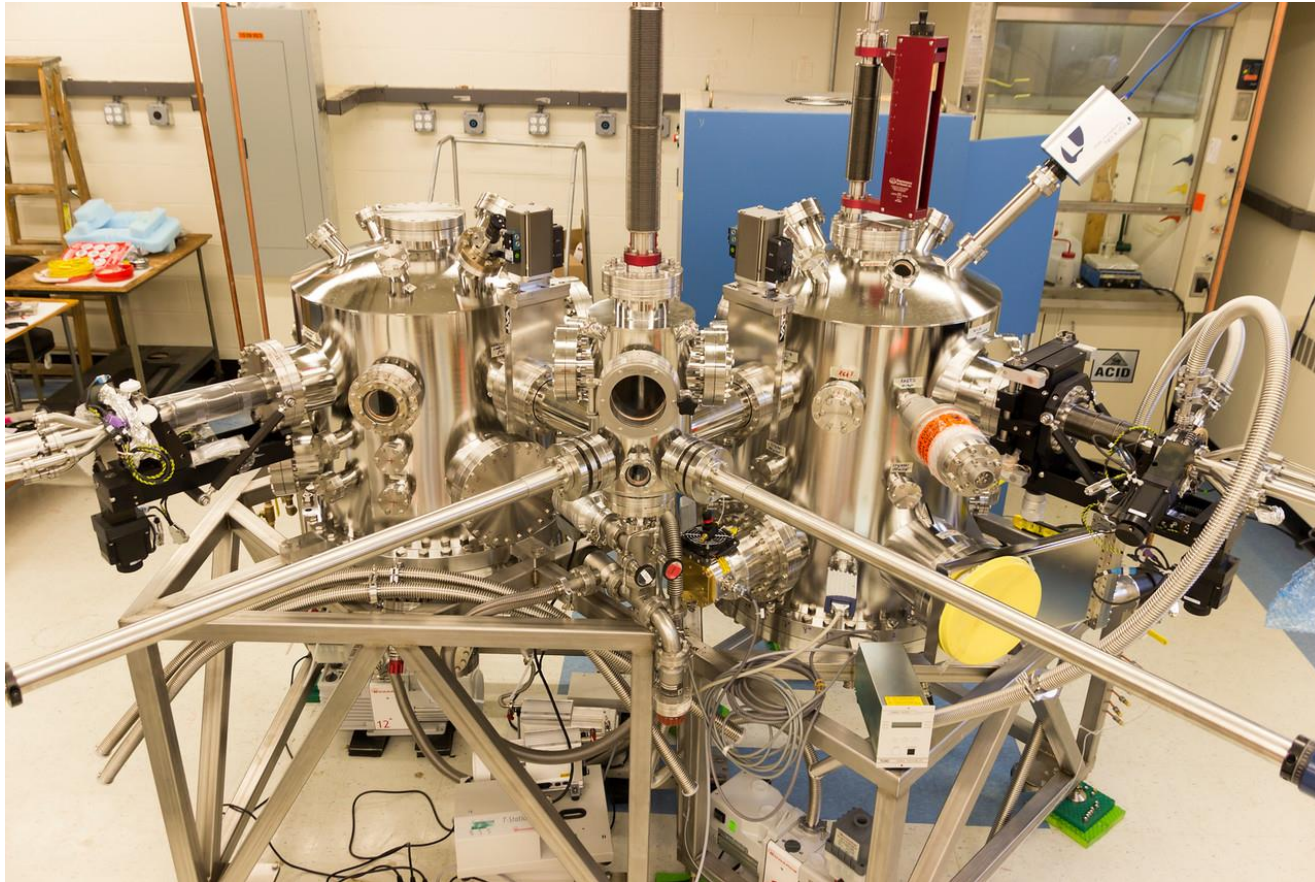
Wire Bonder



➔ Low temp cryostats down to 400 mK, 8 tesla SC magnet and electronics for transport measurements

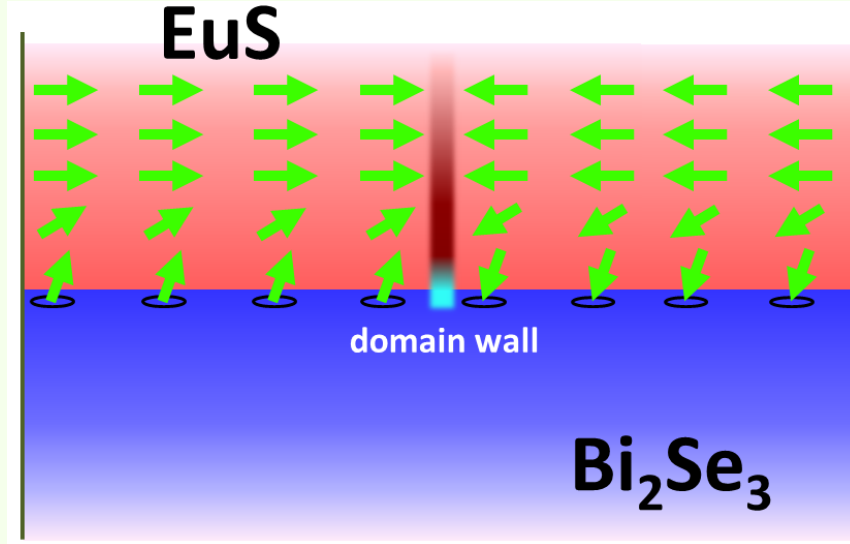
➔ { STM / c- AFM low temp probe, insitu
sample transfer, 280 mK, 5 T in plane H
(April 2015)

New cluster MBE – The Water Molecule !

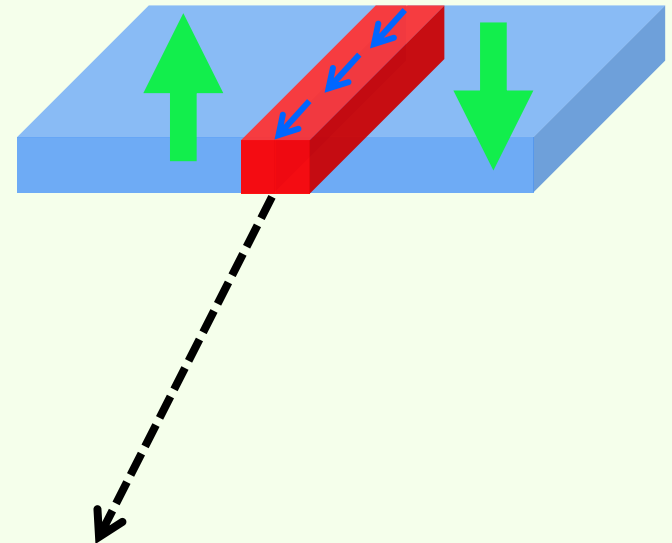


10^{-10} torr, 27 Sources: K-cell, e-Gun evapn, Sputtering
Heating/cooling (100K – 1200K), RHEED (*in situ structure*)
Sample transferable in vacuum (*'vacuum suitcase'*)
Insitu masking, angled evaporation (nano devices)

What happens at the interface:



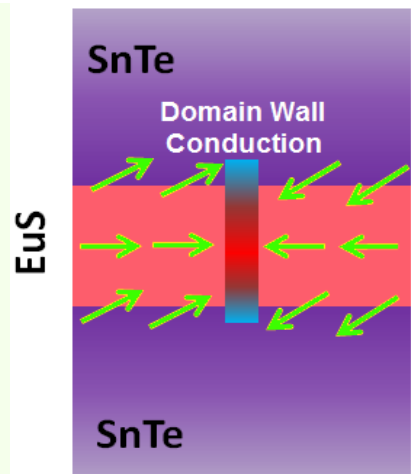
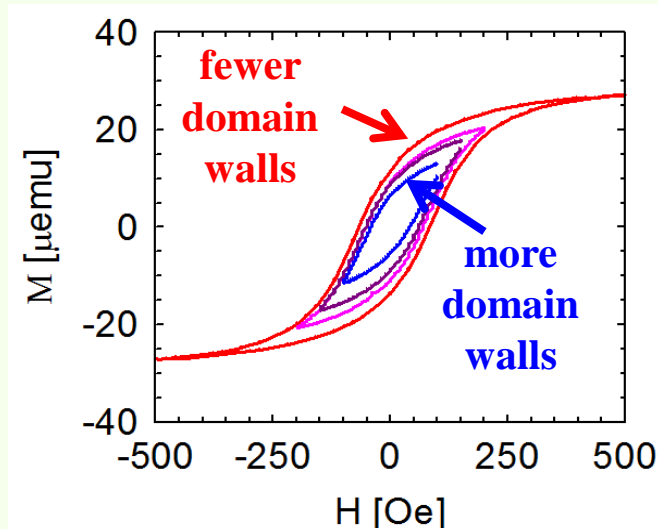
Bi₂Se₃ top surface at coercive field H_c
(R minimum):



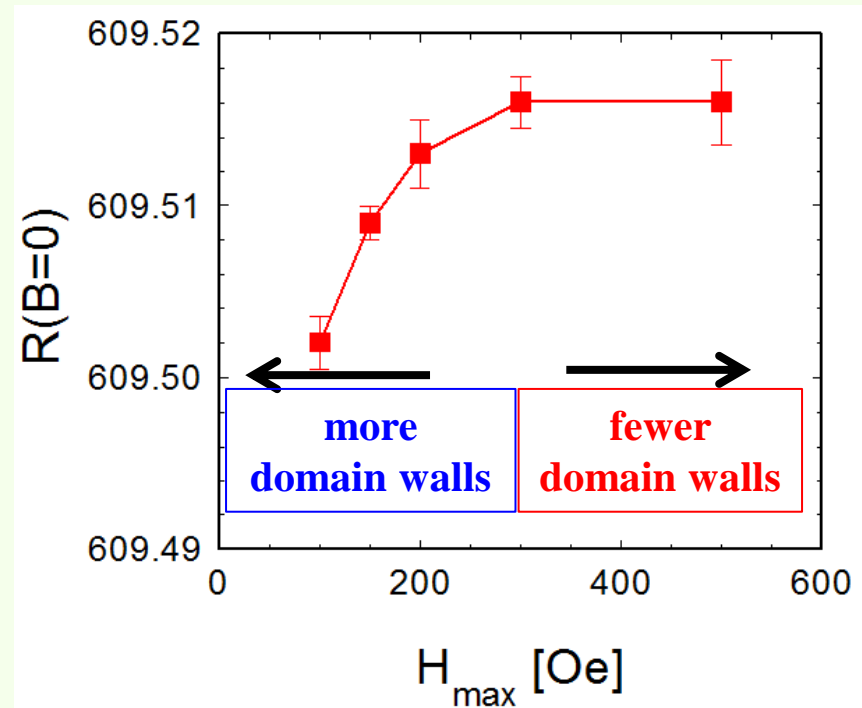
Chiral conduction channel in domain wall (unique in TI)

Proximity induced magnetism in TCI SnTe: Tuning the Number of Domain Walls: the Minor Loop Regime

Minor loops of SnTe-EuS-SnTe Trilayer



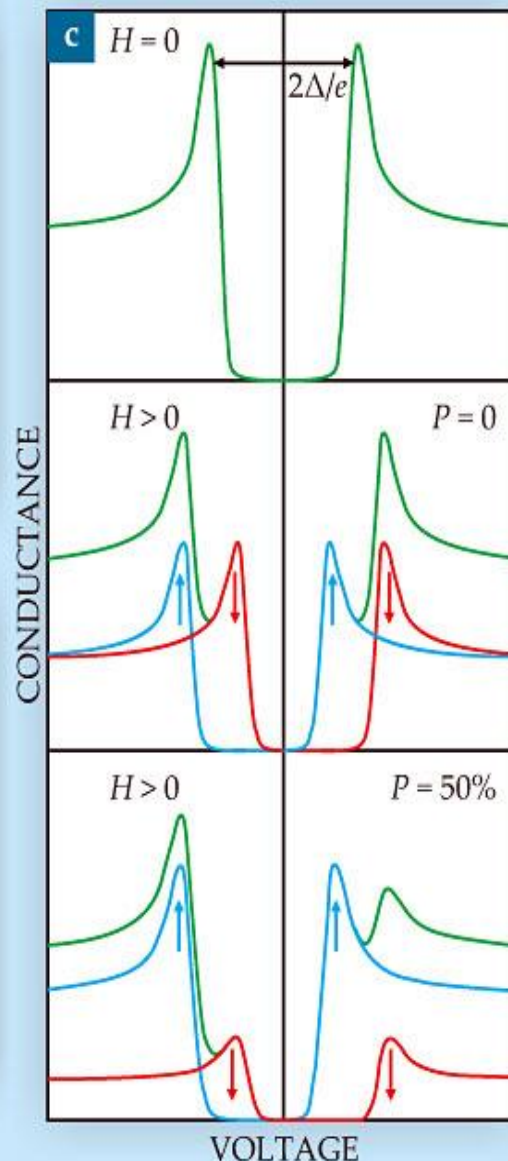
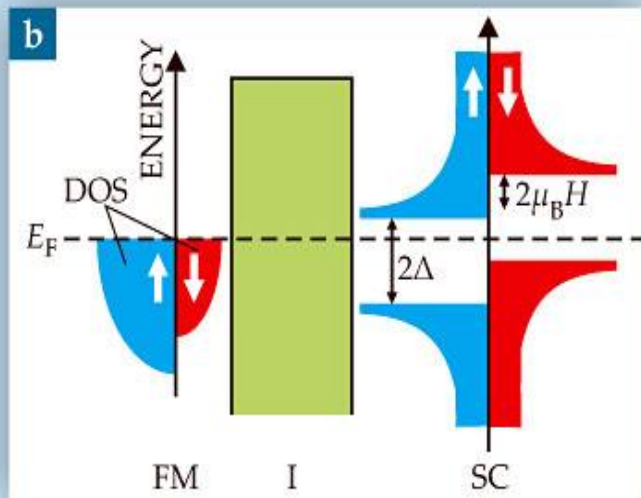
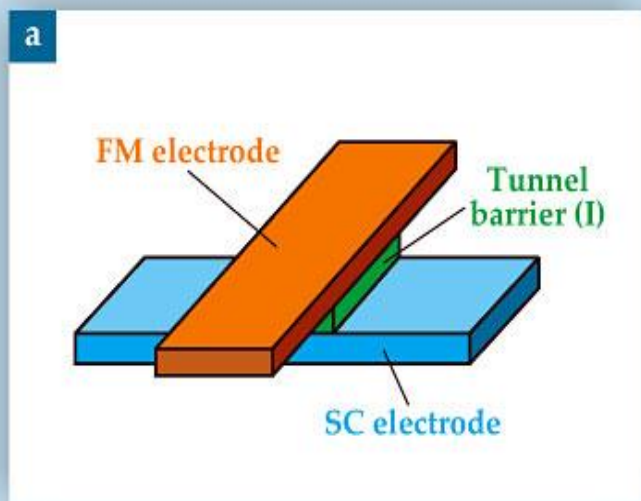
Resistance dependence on domain wall density at $H = 0$



**Magnetic domain walls
enhance sample conduction**

**Correlation between
magnetic domain structure and transport**

Meservey – Tedrow Technique of Spin Polarized Tunneling



**SC / Insulator / Normal metal
tunneling**

Supercond. energy gap

In a magnetic field

Zeeman split states in SC

**SC / Insulator / Normal metal
tunneling**

**SC / Insulator / ferromagnet
tunneling**

Tunneling e spin polarized

$$P = (n_{\uparrow} - n_{\downarrow}) / (n_{\uparrow} + n_{\downarrow})$$

$$P \equiv \frac{(Nv^2)_{\uparrow} - (Nv^2)_{\downarrow}}{(Nv^2)_{\uparrow} + (Nv^2)_{\downarrow}}$$